

*Traveling waves  
and impact-parameter correlations  
in high-energy QCD*

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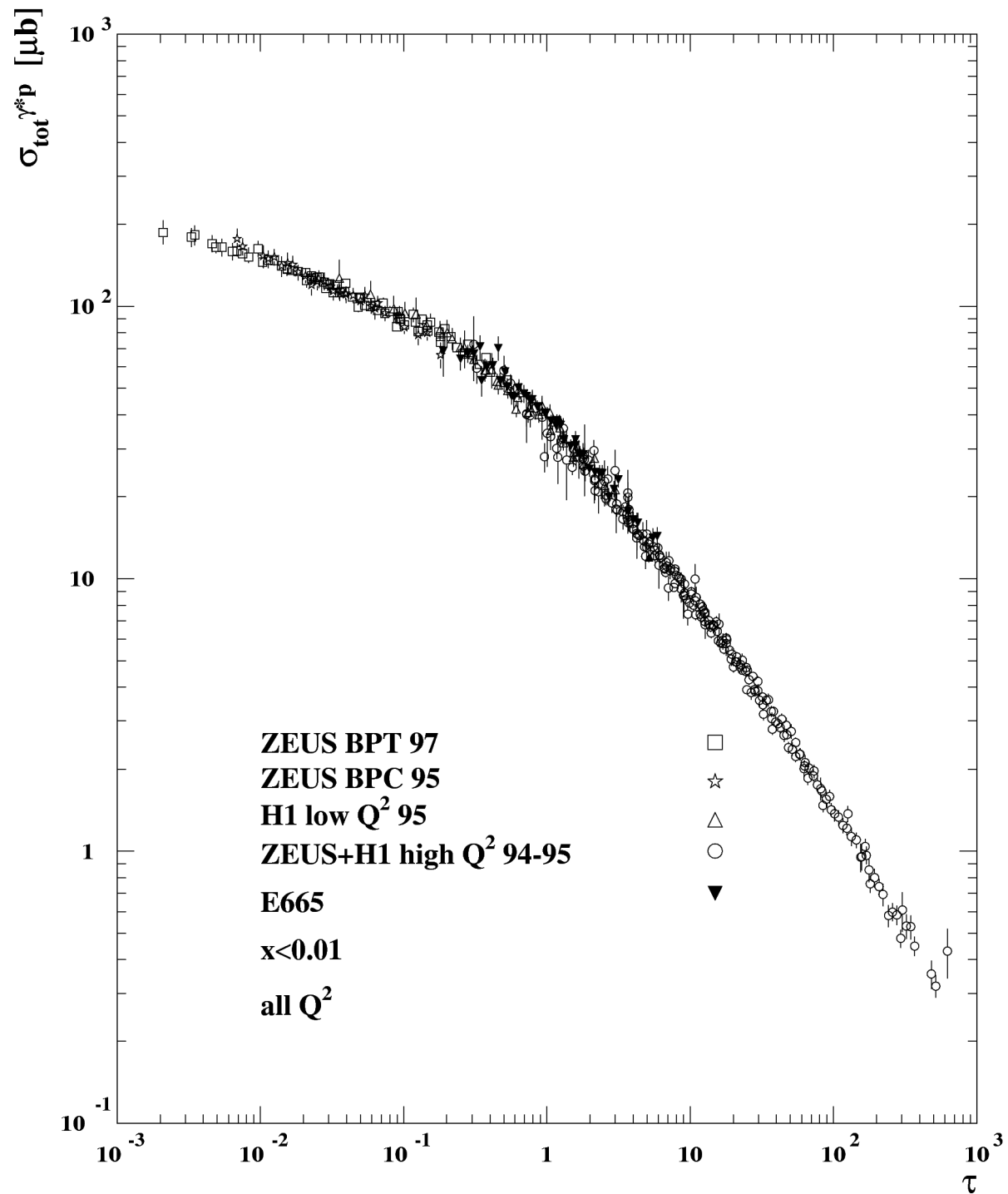


Cracow, January 6, 2009

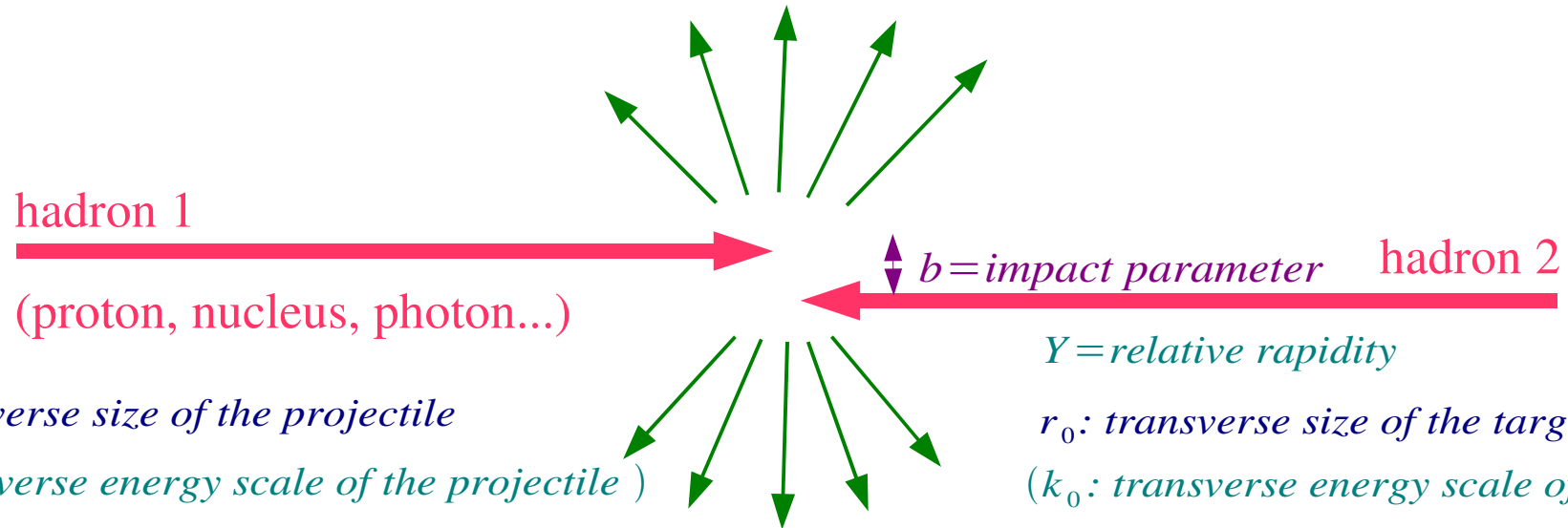
Epiphany conference dedicated to the memory of Jan Kwieciński

# *Geometric scaling*

Staśto,  
Golec-Biernat,  
Kwieciński,  
PRL (2001)



# High energy QCD



$$A(Y, \mathbf{r}) = \int d^2 b A(Y, b, \mathbf{r}) = \text{elastic amplitude}$$

$$A(Y, b, \mathbf{r}) = \text{fixed impact parameter amplitude} \leq 1$$

(High) energy dependence of QCD amplitudes?

# QCD and statistical physics

Iancu, Mueller, Munier (2004)

*We identified the universality class of high-energy scattering at fixed impact parameter as that of **one-dimensional** (« space » variable:  $\log k^2$  ) **reaction-diffusion processes** evolving in « time »  $Y$*

- ★ simple picture of high energy scattering, based on the parton model
- ★ connects the QCD problem to more general physics and mathematics
- ★ new results for QCD amplitudes!

# QCD and statistical physics

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- ★ connects the QCD problem to more general physics and mathematics
- ★ new results for QCD amplitudes!

**But: this identification is still a conjecture!**

In particular, it is not completely clear that QCD may be reduced to a one-dimensional problem

$$A(Y, \cancel{b}, r)$$

**Condition: each impact parameter evolves independently.**

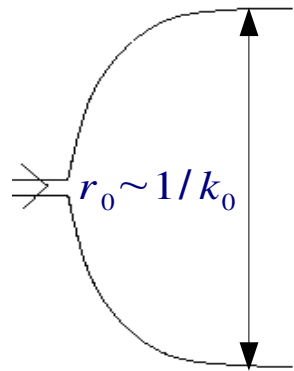
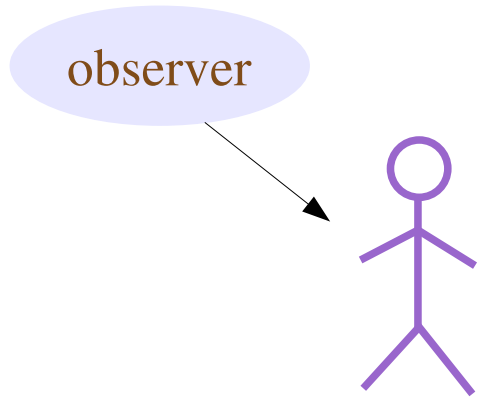
***This talk: a numerical check of this statement***

Munier, Salam, Soyez (2008)

# *Outline*

- ★ High energy QCD and one-dimensional stochastic processes
- ★ Independence of different impact parameters in a toy model

# *How a high rapidity hadron looks*

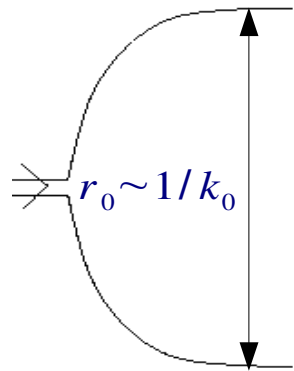
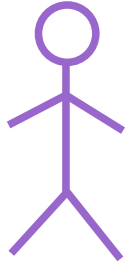


$$Y_0 = 0$$

rapidity in the frame  
of the observer

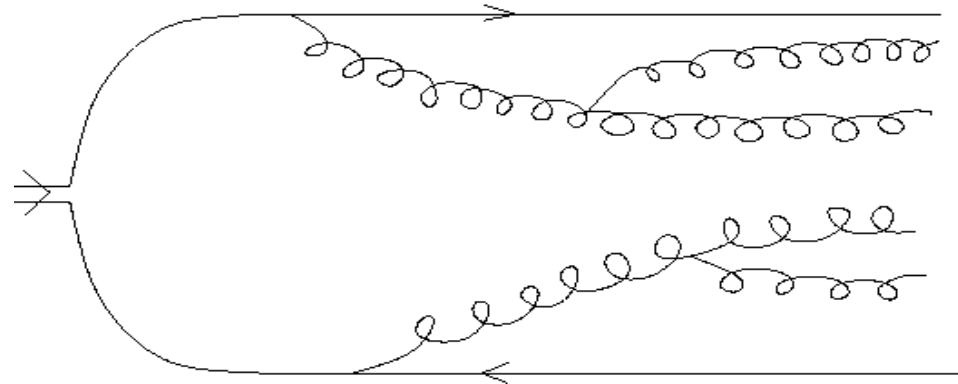
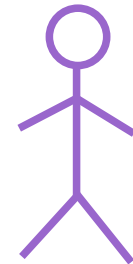
# How a high rapidity hadron looks

observer



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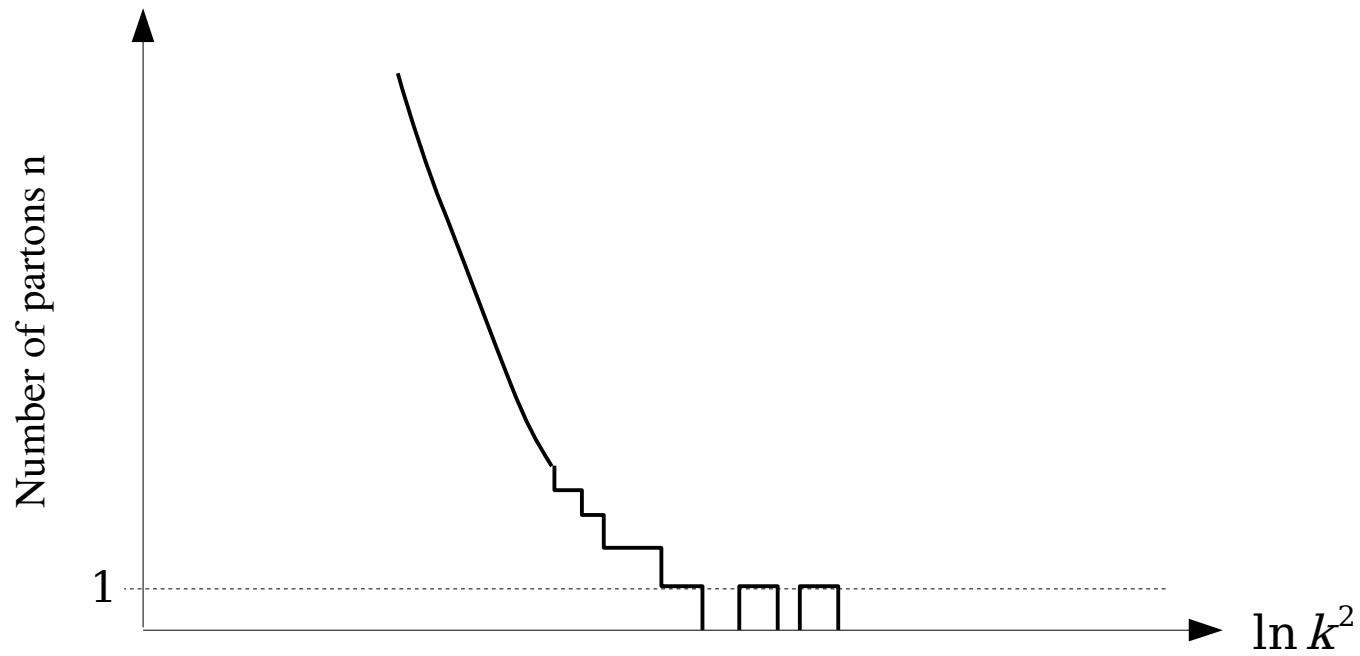
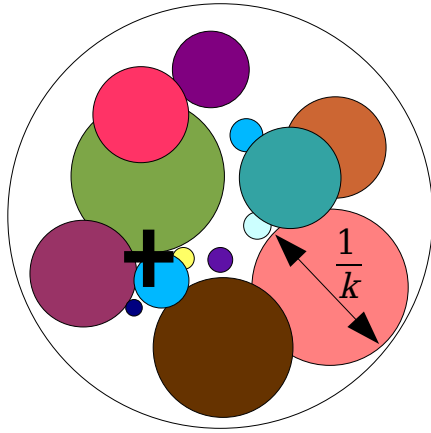
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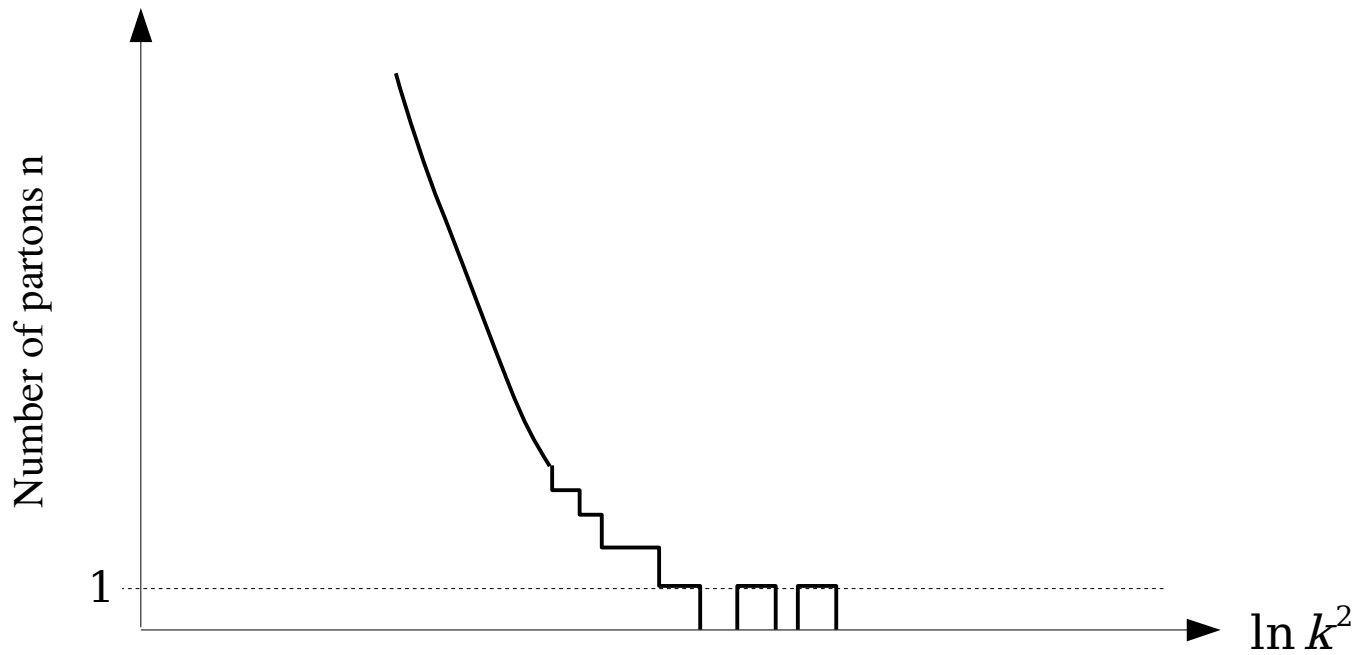
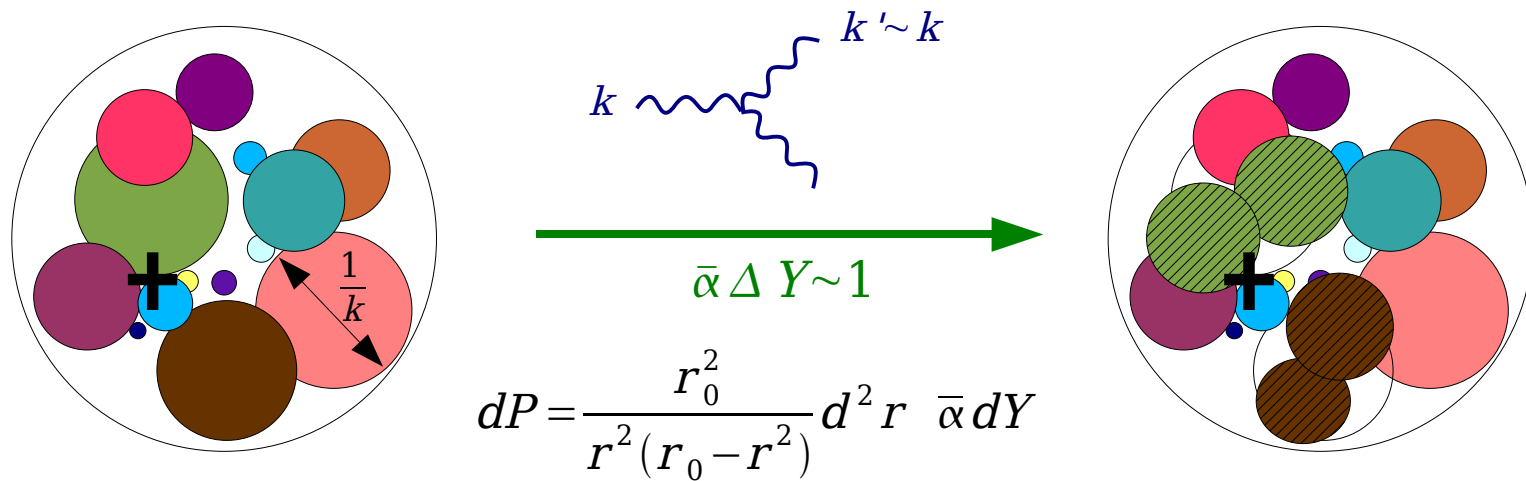
$$Y_1 > Y_0$$



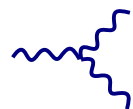
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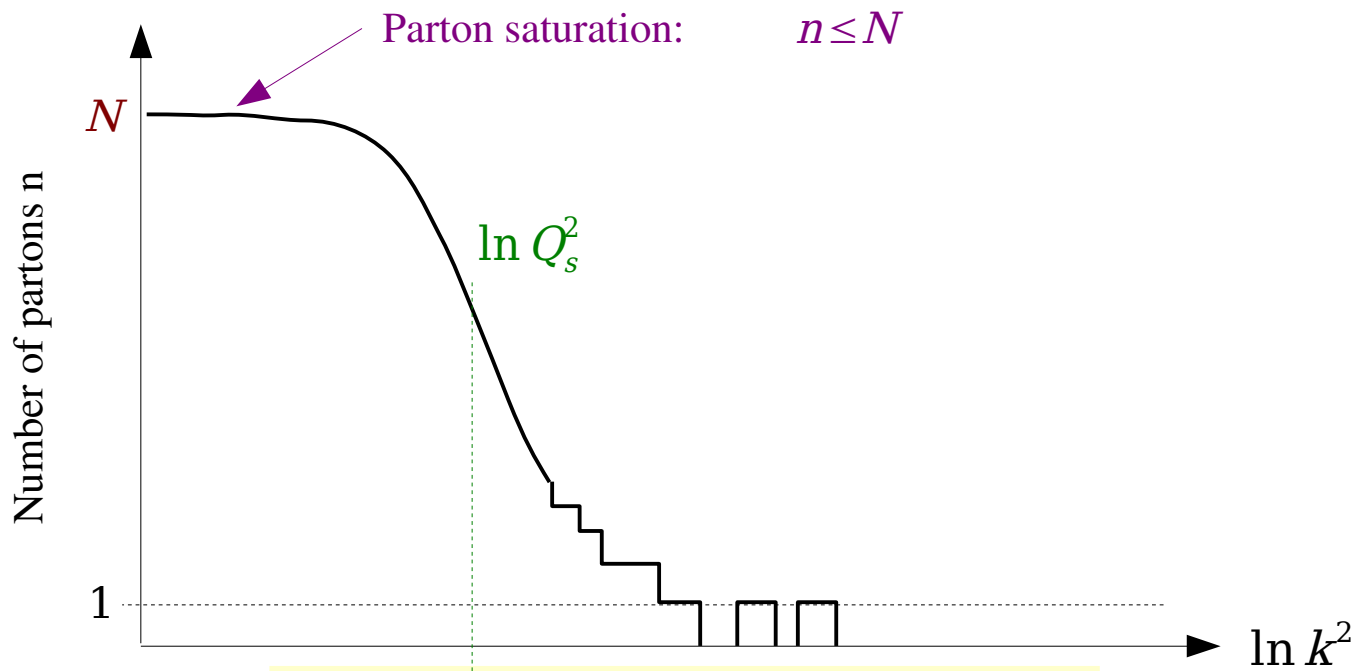
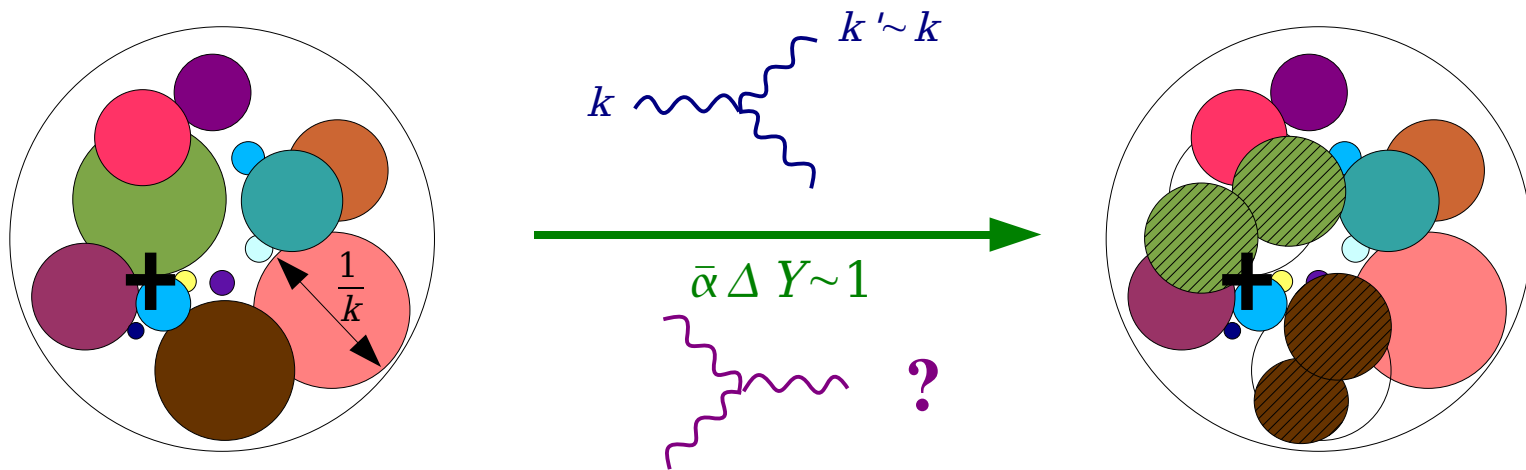
$$\partial_{\bar{\alpha} Y} n = \chi(-\partial_{\ln k^2}) n + \sqrt{n} \nu$$



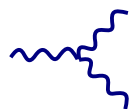
$BFKL \sim \partial_{\ln k^2}^2 n + n$

Noise term due to discreteness

# How a high rapidity hadron looks



$$\partial_{\bar{\alpha} Y} n = \chi(-\partial_{\ln k^2}) n - \frac{n^2}{N} + \sqrt{n} \nu$$

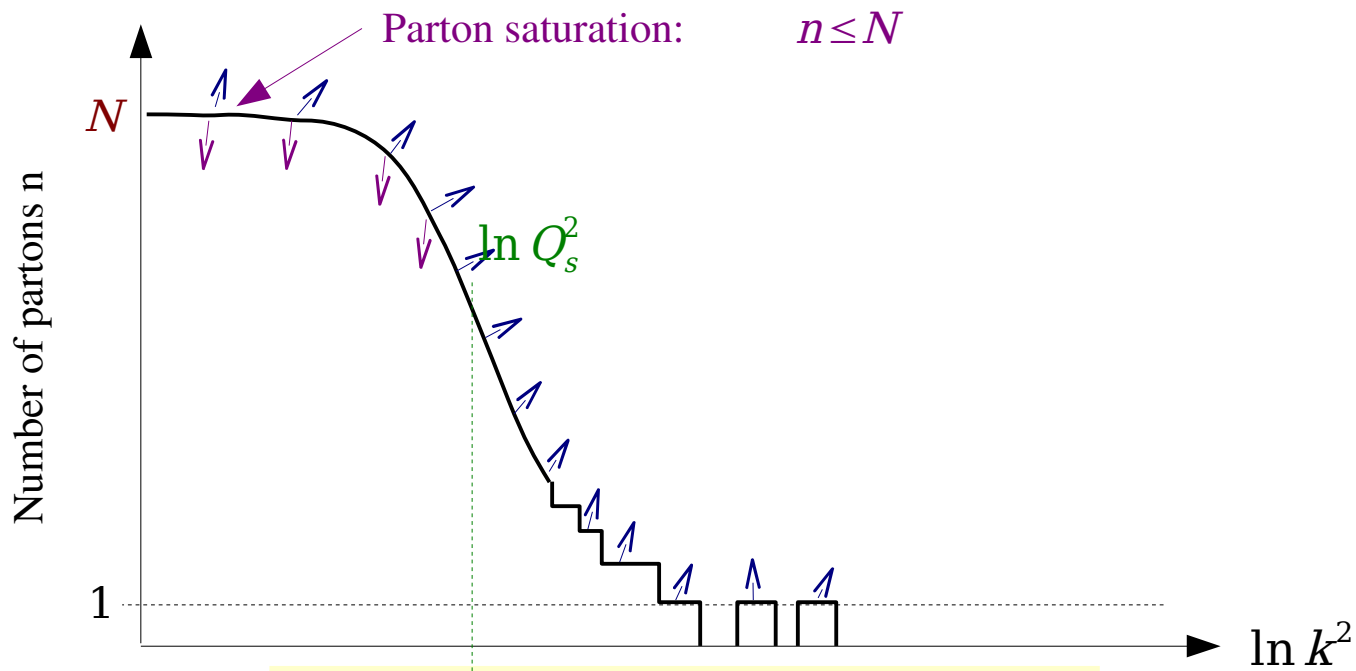
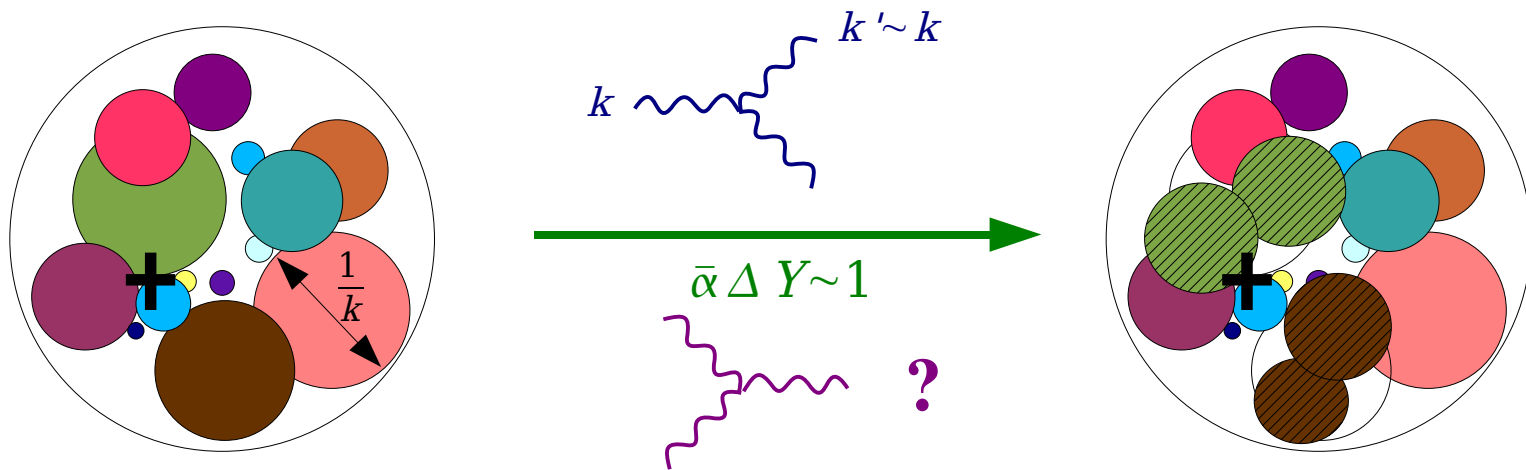


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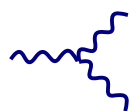


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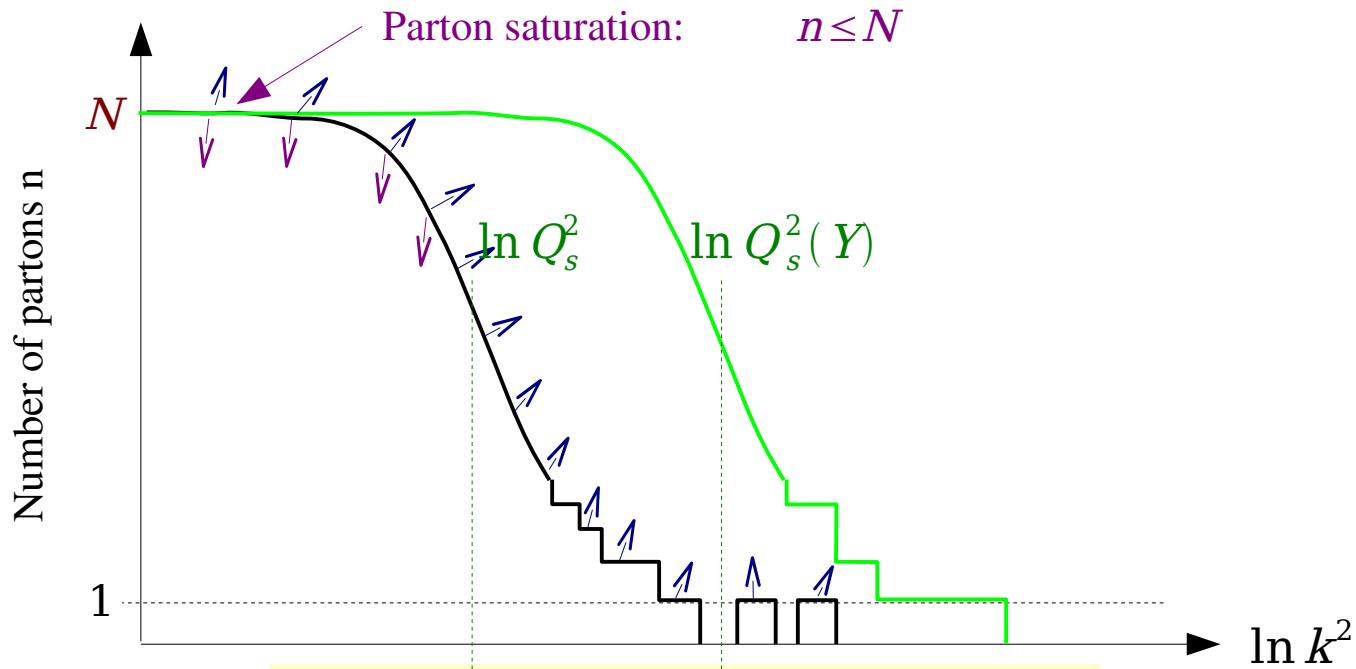
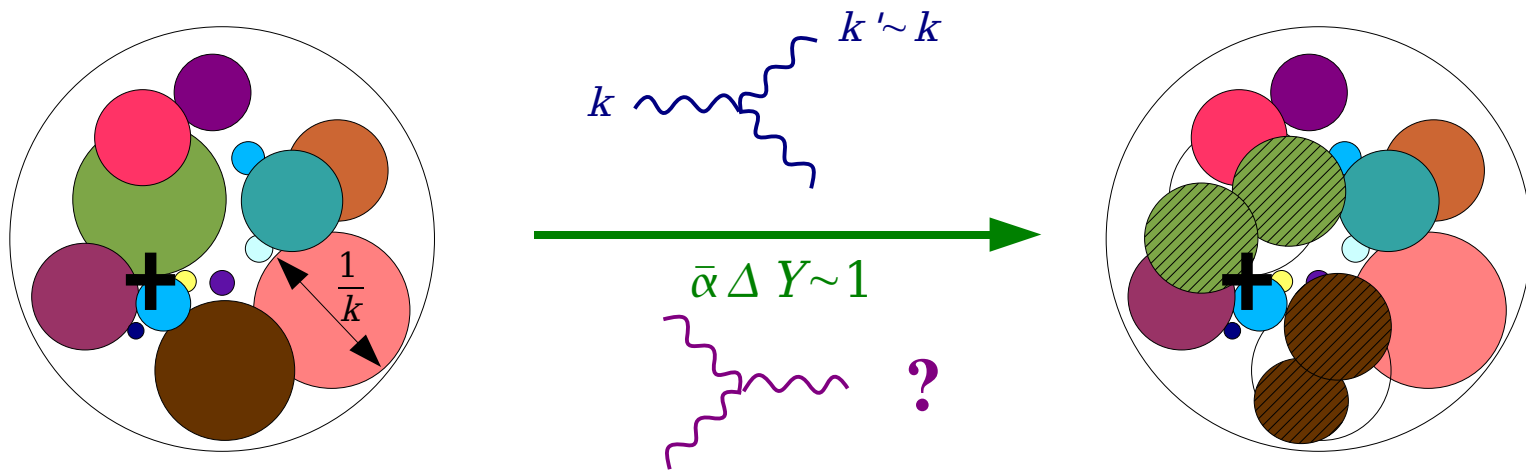


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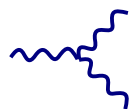


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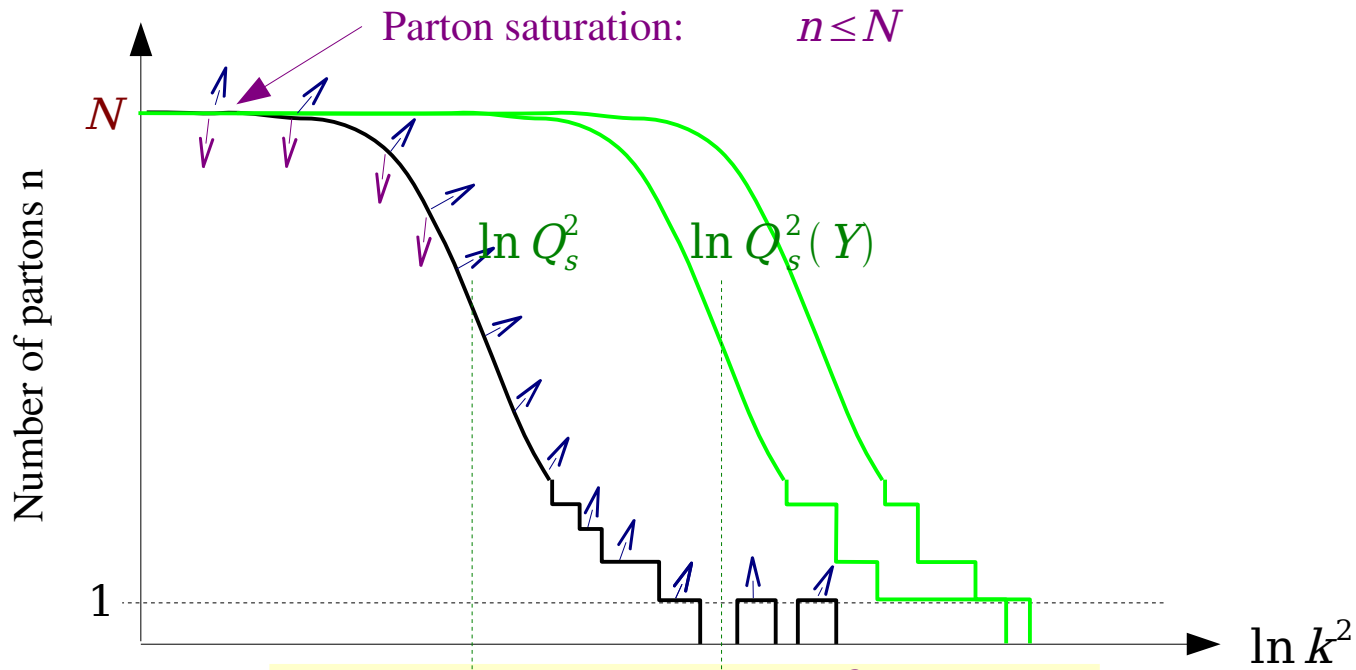
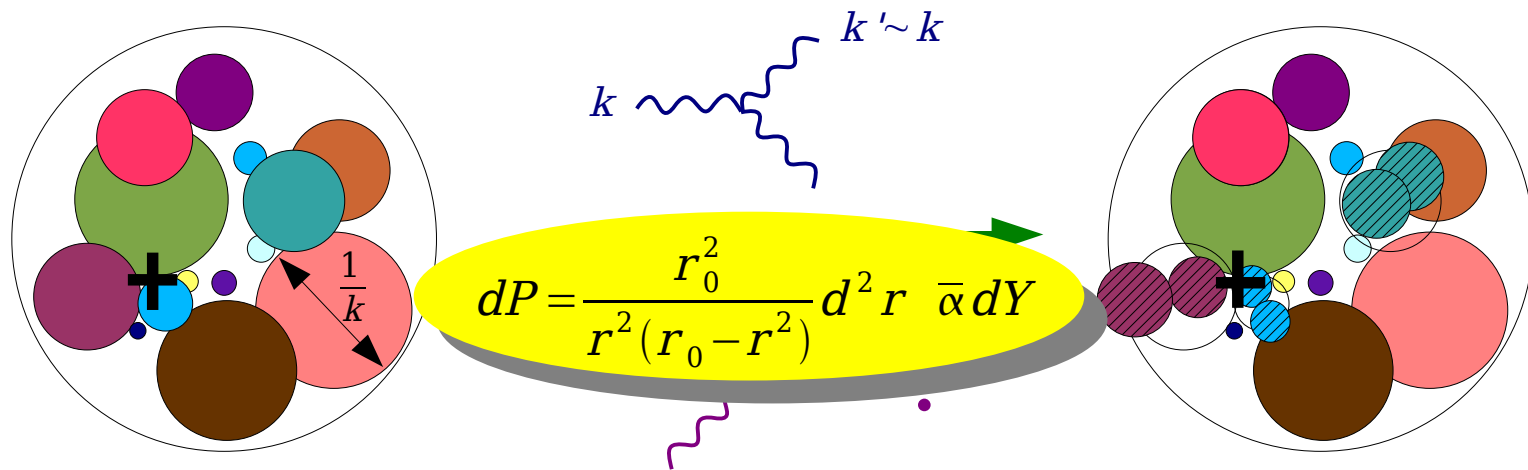


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Noise term due to discreteness

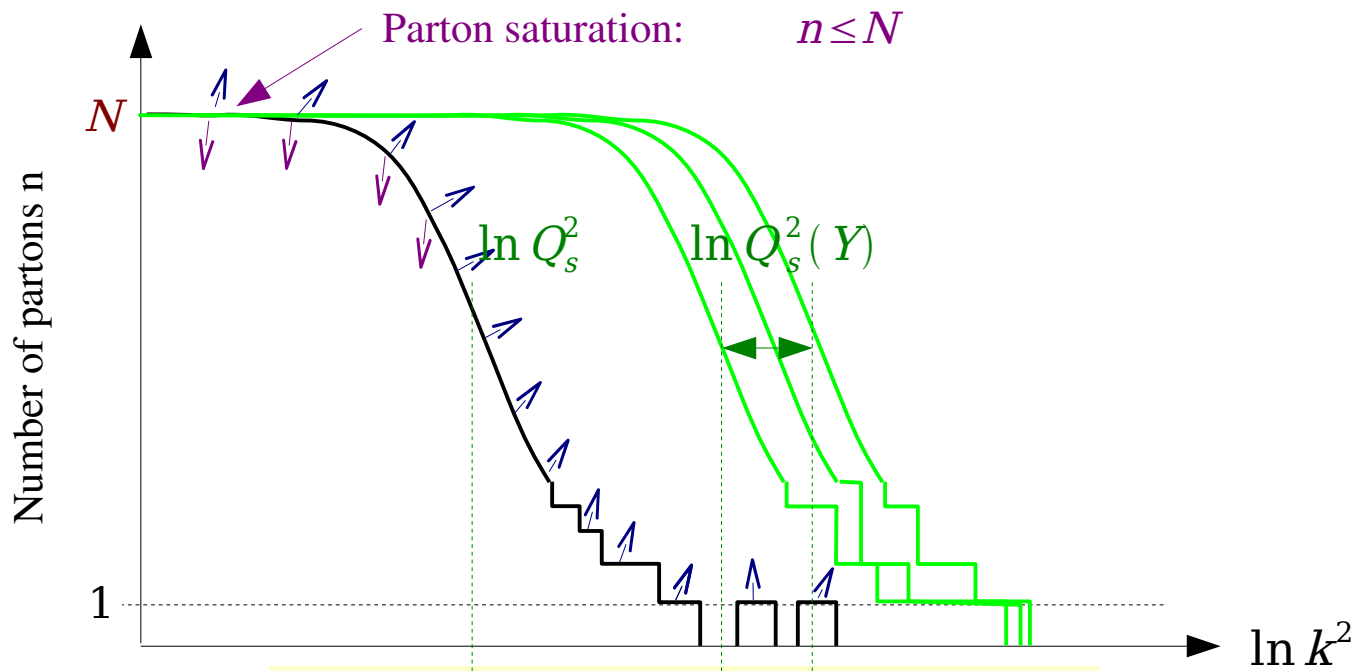
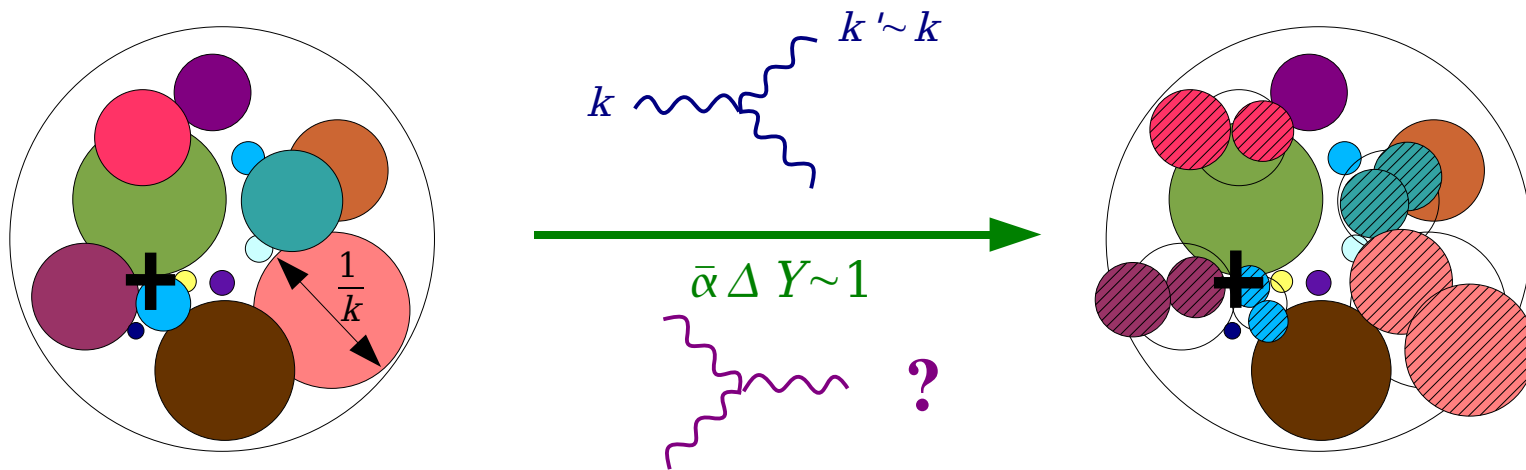
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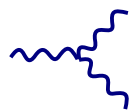
$\partial_{\bar{\alpha} Y} n = \chi(-\partial_{\ln k^2}) n - \frac{n^2}{N} + \sqrt{n} \nu$

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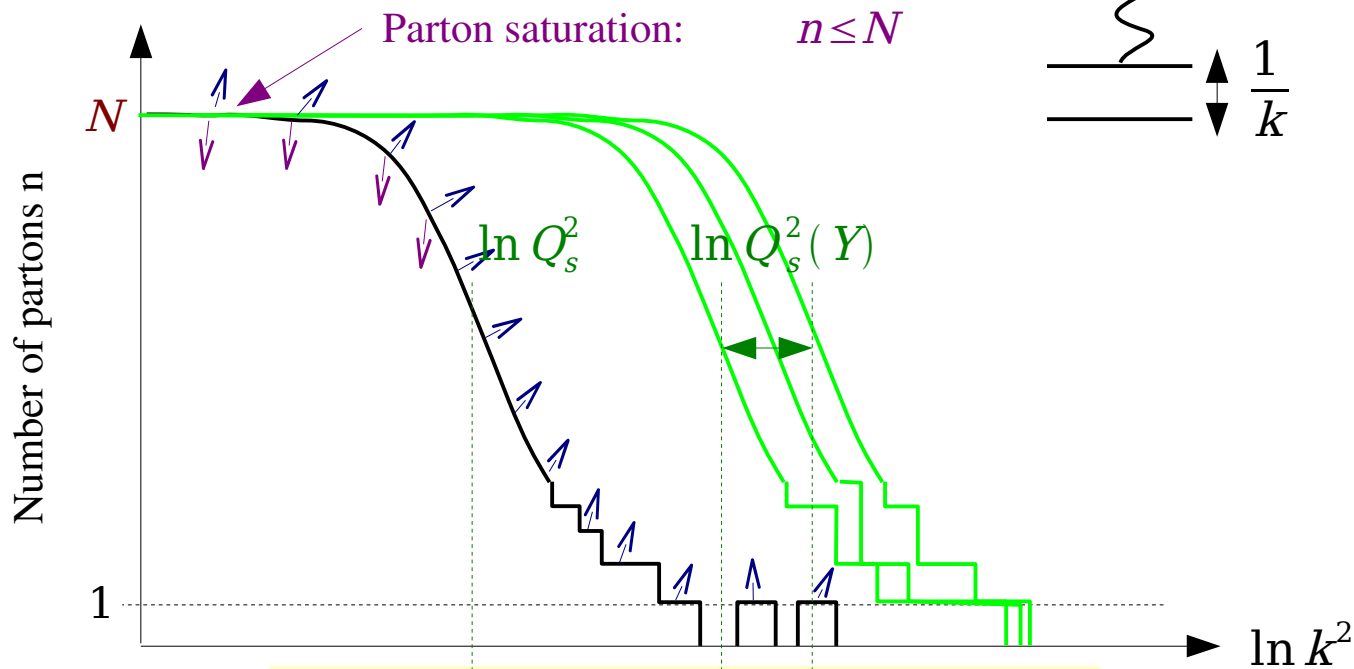
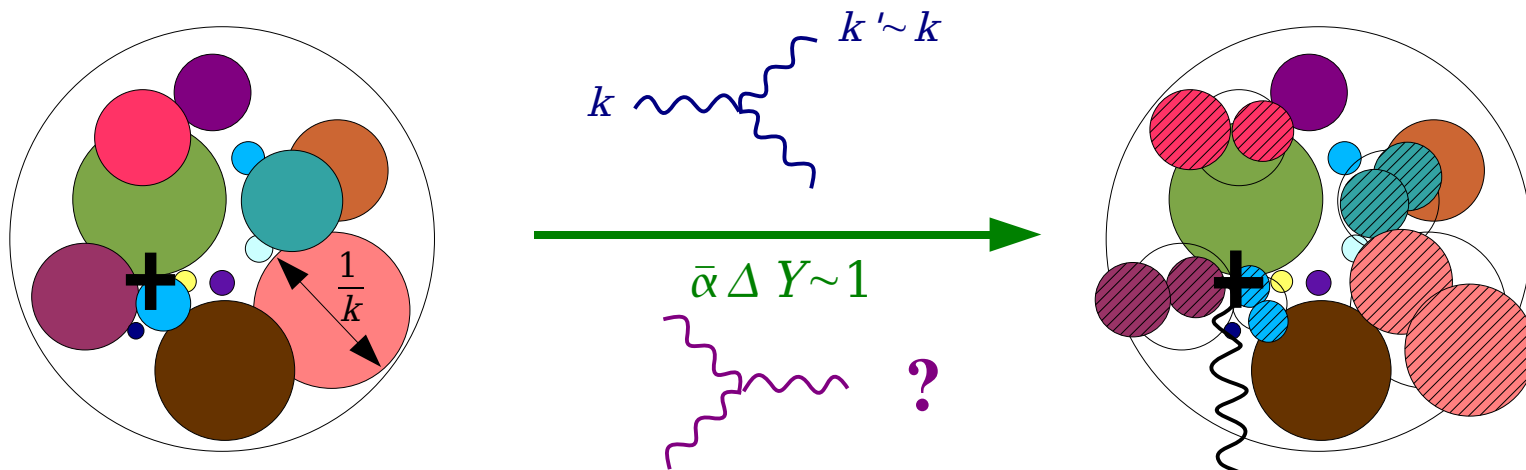


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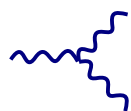
Noise term due to discreteness

# How a high rapidity hadron looks



$$T(k) \sim \alpha_s^2 n(k)$$

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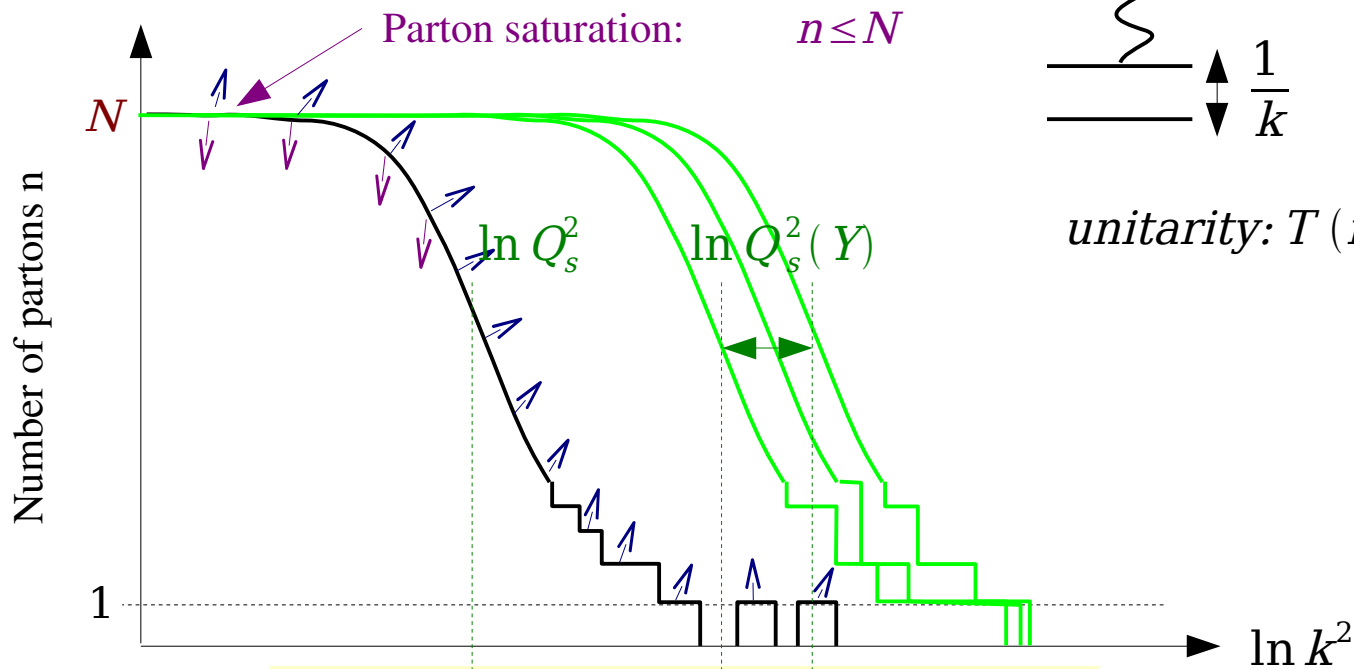
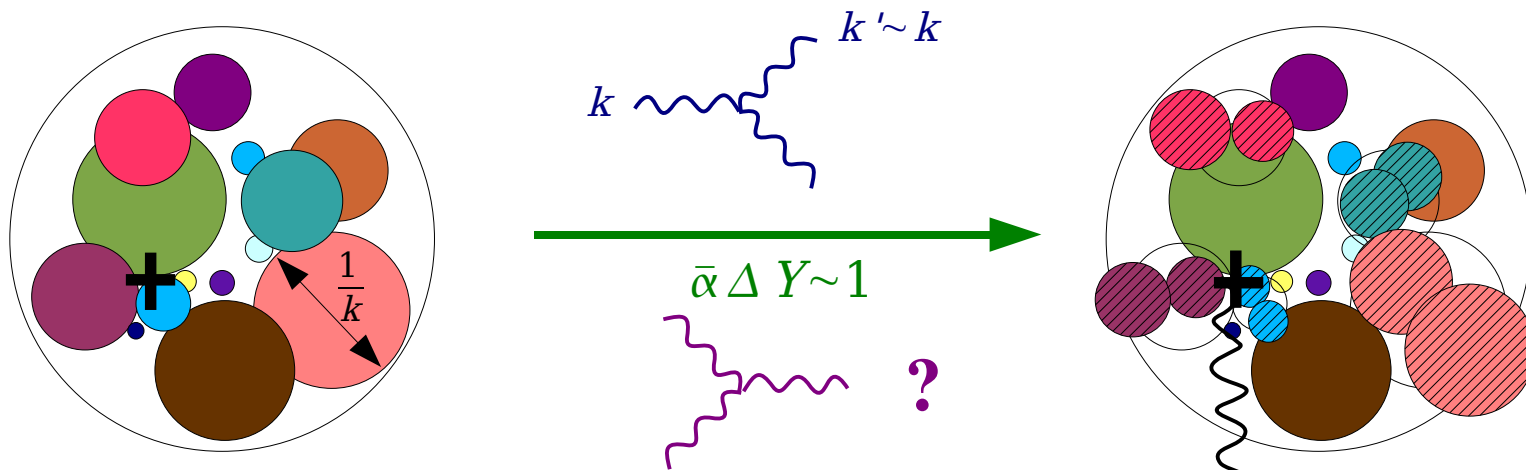


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Noise term due to discreteness



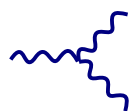
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$$\text{unitarity: } T(r) \leq 1 \Rightarrow N = \frac{1}{\alpha_s^2}$$

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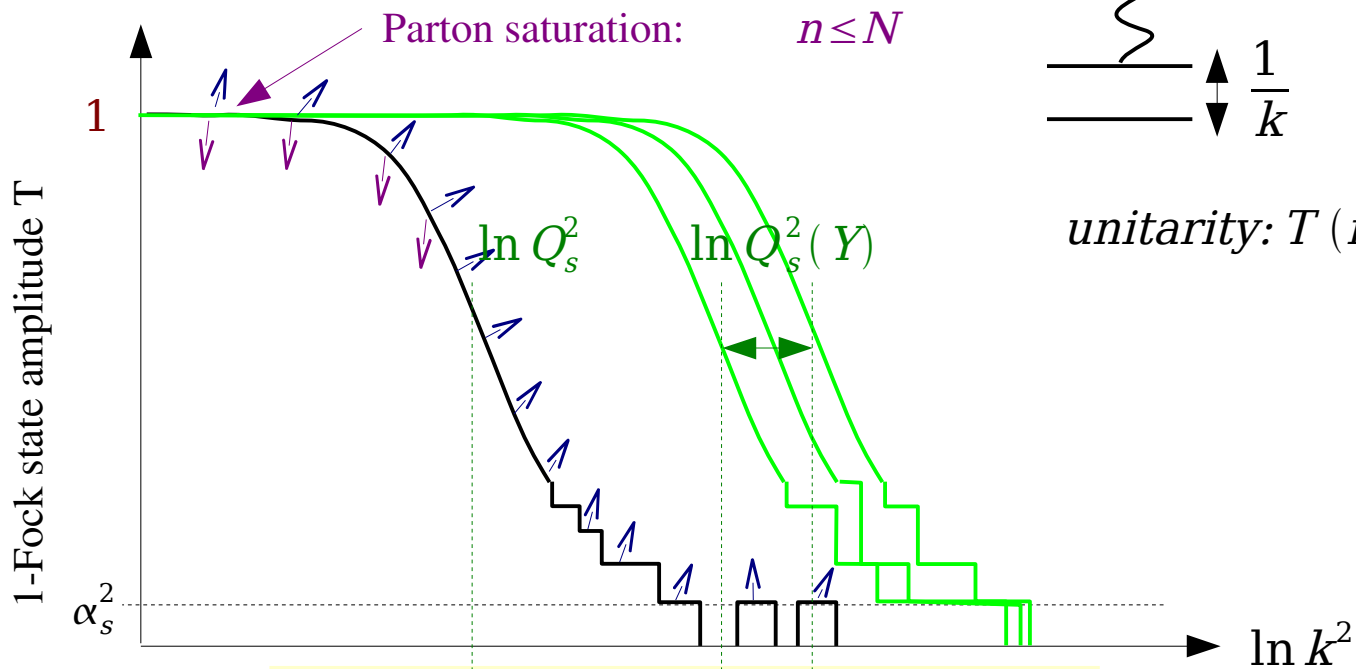
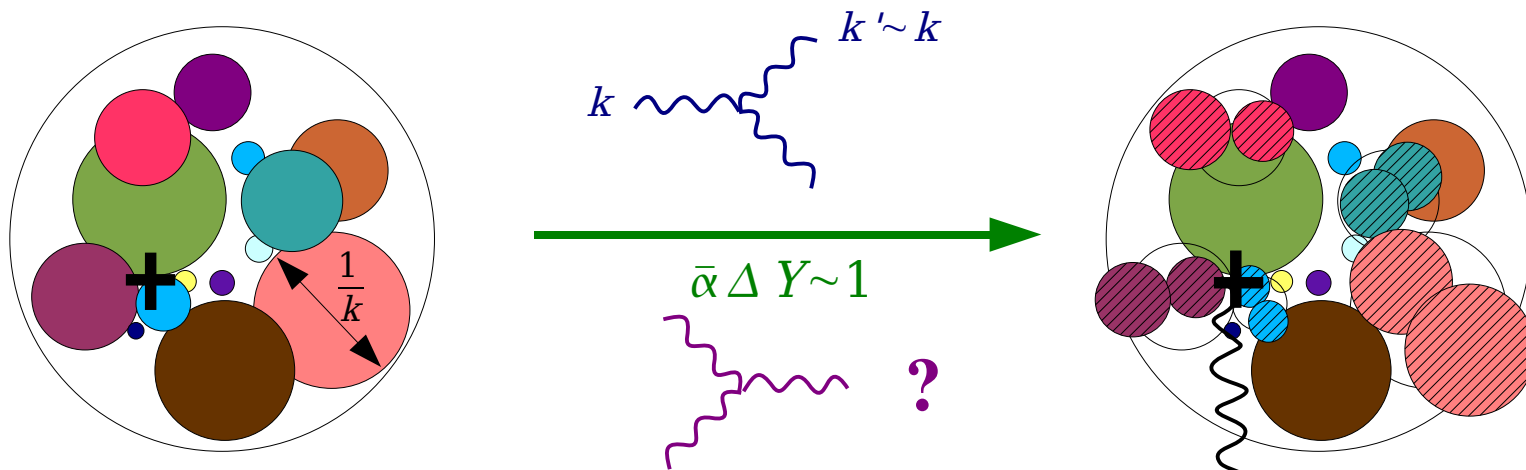


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Noise term due to discreteness

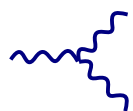
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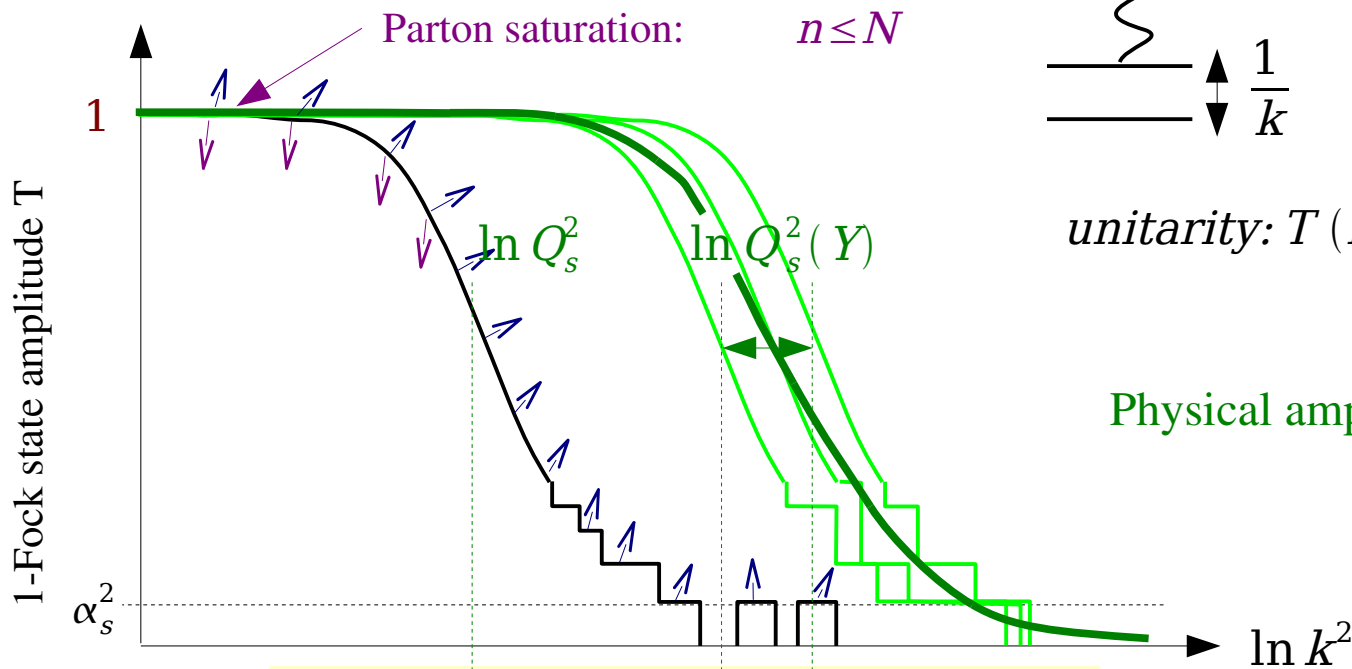
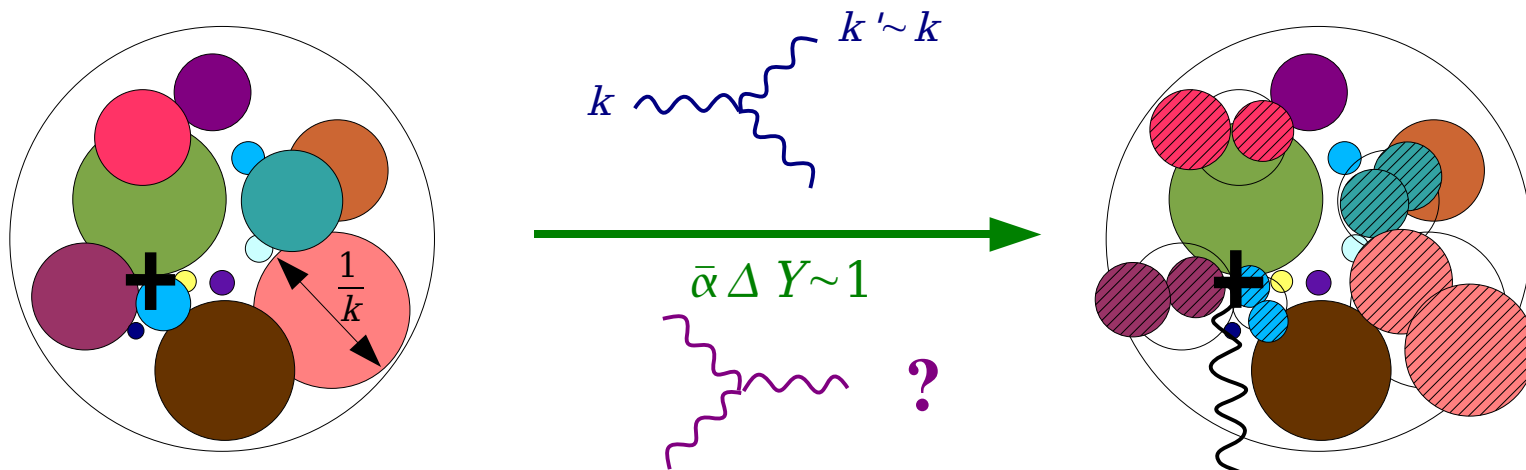


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Noise term due to discreteness

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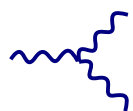


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$$\text{unitarity: } T(r) \leq 1 \Rightarrow N = \frac{1}{\alpha_s^2}$$

$$\text{Physical amplitude: } A = \langle T \rangle$$

$$\partial_{\bar{\alpha} Y} T = \chi(-\partial_{\ln k^2}) T - T^2 + \alpha_s \sqrt{T} \nu$$



$$BFKL \sim \partial_{\ln k^2}^2 T + T$$



Noise term due to discreteness

# QCD and reaction-diffusion

$$\partial_{\bar{\alpha} Y} T = \chi(-\partial_{\ln k^2}) T - T^2 + \alpha_s \sqrt{T} \nu$$

Similar to the sFKPP equation  $\partial_t T = \partial_x^2 T + T - T^2 + \sqrt{\frac{2}{N}} T \nu$

Fisher; Kolmogorov,  
Petrovsky, Piscunov (1937)

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Fisher; Kolmogorov,  
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## Predictions for QCD amplitudes

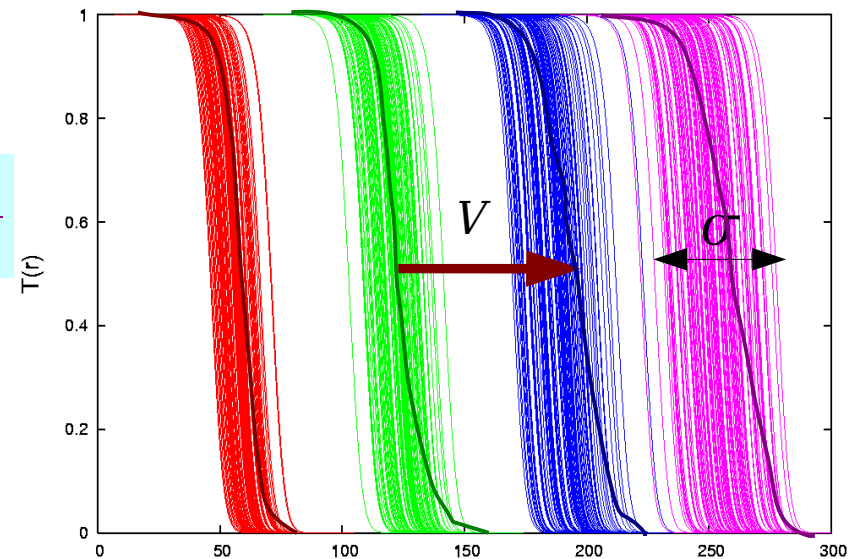
Shape of the *partonic* amplitude:  $T \sim (rQ_s(Y))^{2\gamma_0}$

Saturation scale:

$$V = \frac{d}{d(\bar{\alpha} Y)} \langle \ln Q_s^2 \rangle = \frac{\chi(\gamma_0)}{\gamma_0} - \frac{\pi^2 \gamma_0 \chi''(\gamma_0)}{2 \ln^2(1/\alpha_s^2)} + \pi^2 \gamma_0^2 \chi''(\gamma_0) \frac{3 \ln \ln(1/\alpha_s^2)}{\gamma_0 \ln^3(1/\alpha_s^2)}$$

$$\langle \ln^n Q_s^2 \rangle_{\text{cumulant}} = \pi^2 \gamma_0^2 \chi''(\gamma_0) \frac{n! \zeta(n)}{\gamma_0^n} \left[ \frac{\bar{\alpha} Y}{\ln^3(1/\alpha_s^2)} \right]$$

$$\Rightarrow A \sim A \left( \frac{r^2 Q_s^2(Y)}{\sqrt{\frac{\bar{\alpha} Y}{\ln^3(1/\alpha_s^2)}}} \right)$$

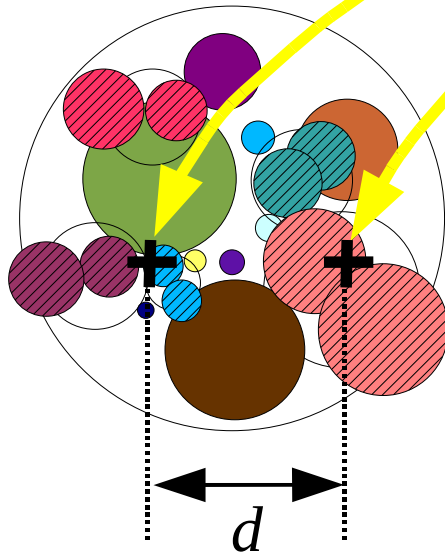


*Traveling waves*

Brunet, Derrida, Mueller, Munier (2004)

These formulas are independent of the precise form of the stochasticity and of the nonlinearity.

# *Independence of the different impact parameters?*



The amplitudes at **this** and **that** impact parameters are independent as soon as

$$1/\langle Q_s(Y, b_1) \rangle, 1/\langle Q_s(Y, b_2) \rangle < d$$

As soon as the distance between the probed impact parameters is larger than the relevant distance scale of the evolution (=the inverse saturation scale), the amplitudes measured at the two impact parameters should be independent.

*Supported by (too) simple analytical estimates  
(fluctuations neglected...)*

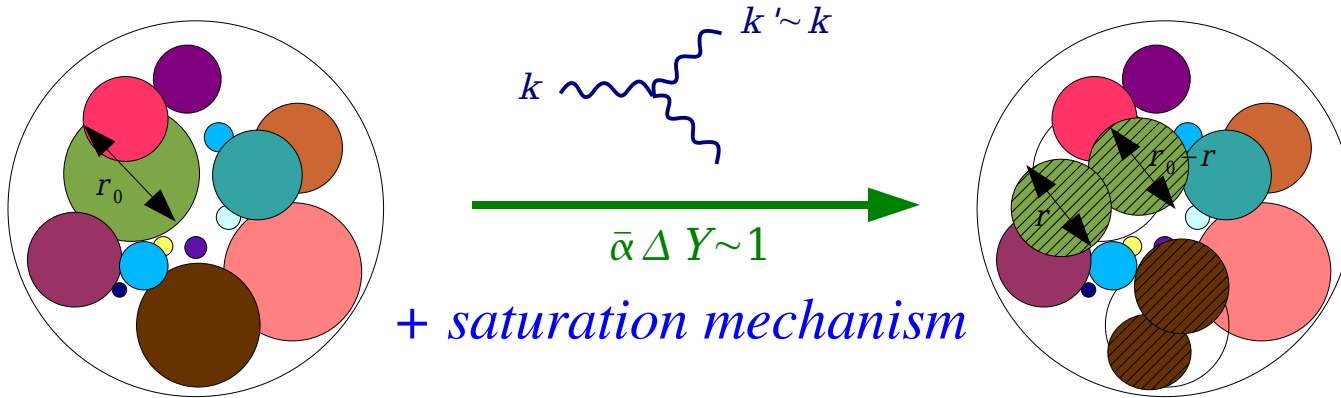
# *Outline*

- ★ High energy QCD and one-dimensional stochastic processes
- ★ Independence of different impact parameters in a toy model

# Toy model with impact-parameter dependence

QCD:

$$dP = \frac{r_0^2}{r^2(r_0 - r^2)} d^2 r \bar{\alpha} dY$$



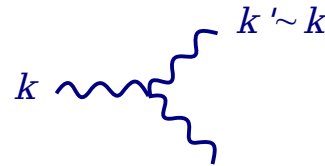
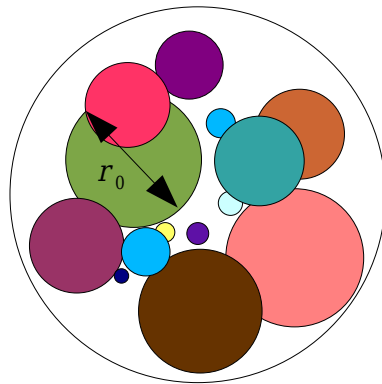
***Too complicated!***



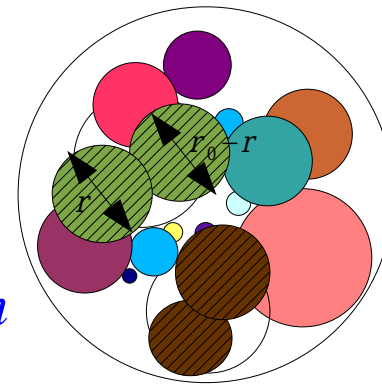
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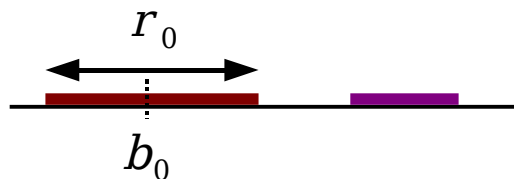
+ saturation mechanism



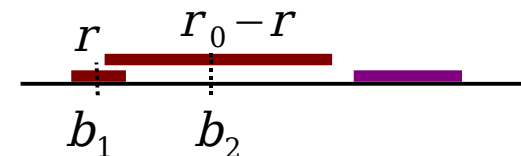
**Too complicated!**

Toy model:

$$dP = \frac{|r_0|}{|r||r_0 - r|} dr dY$$



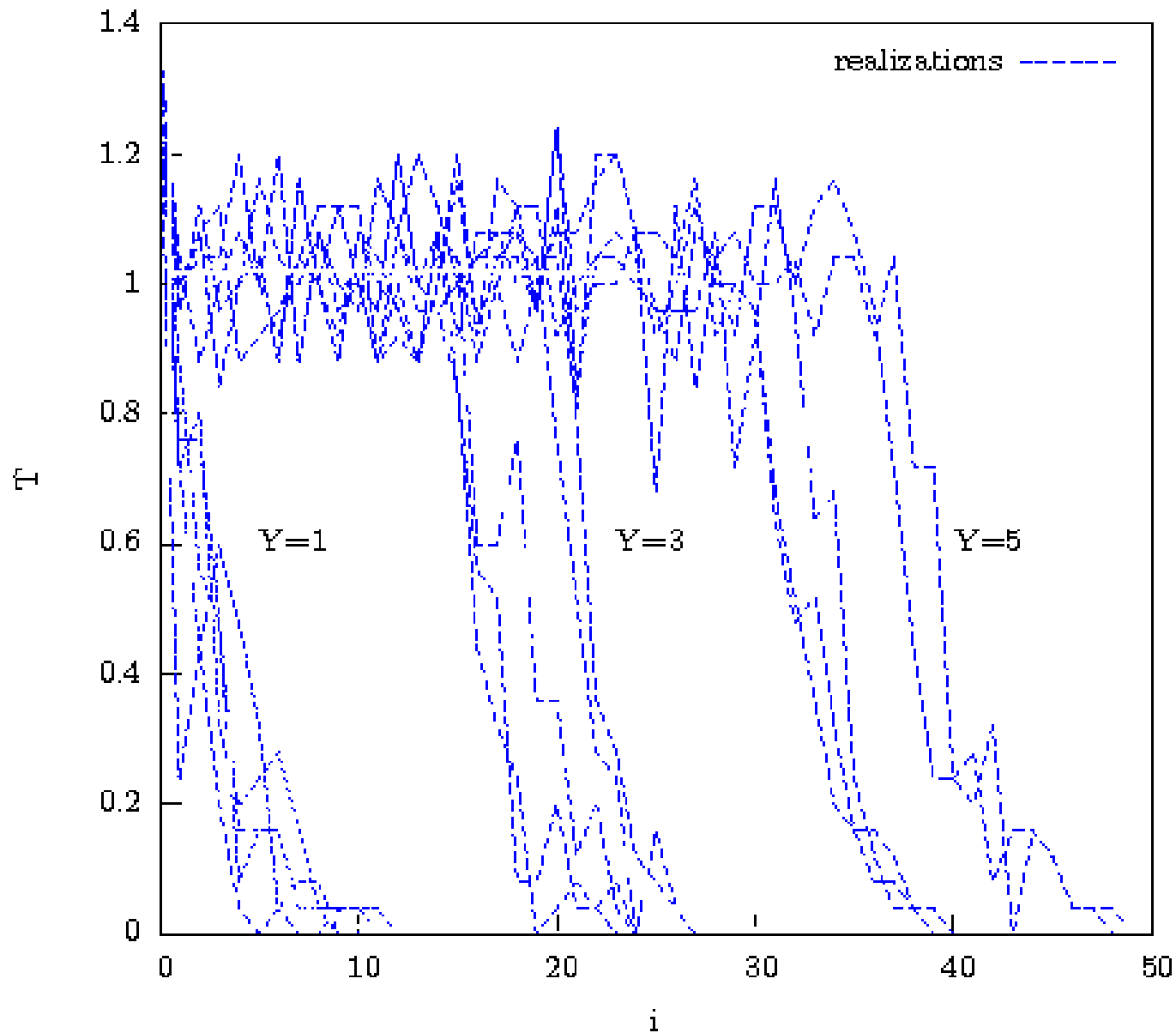
« interval splitting »



+ discretized sizes + saturation mechanism

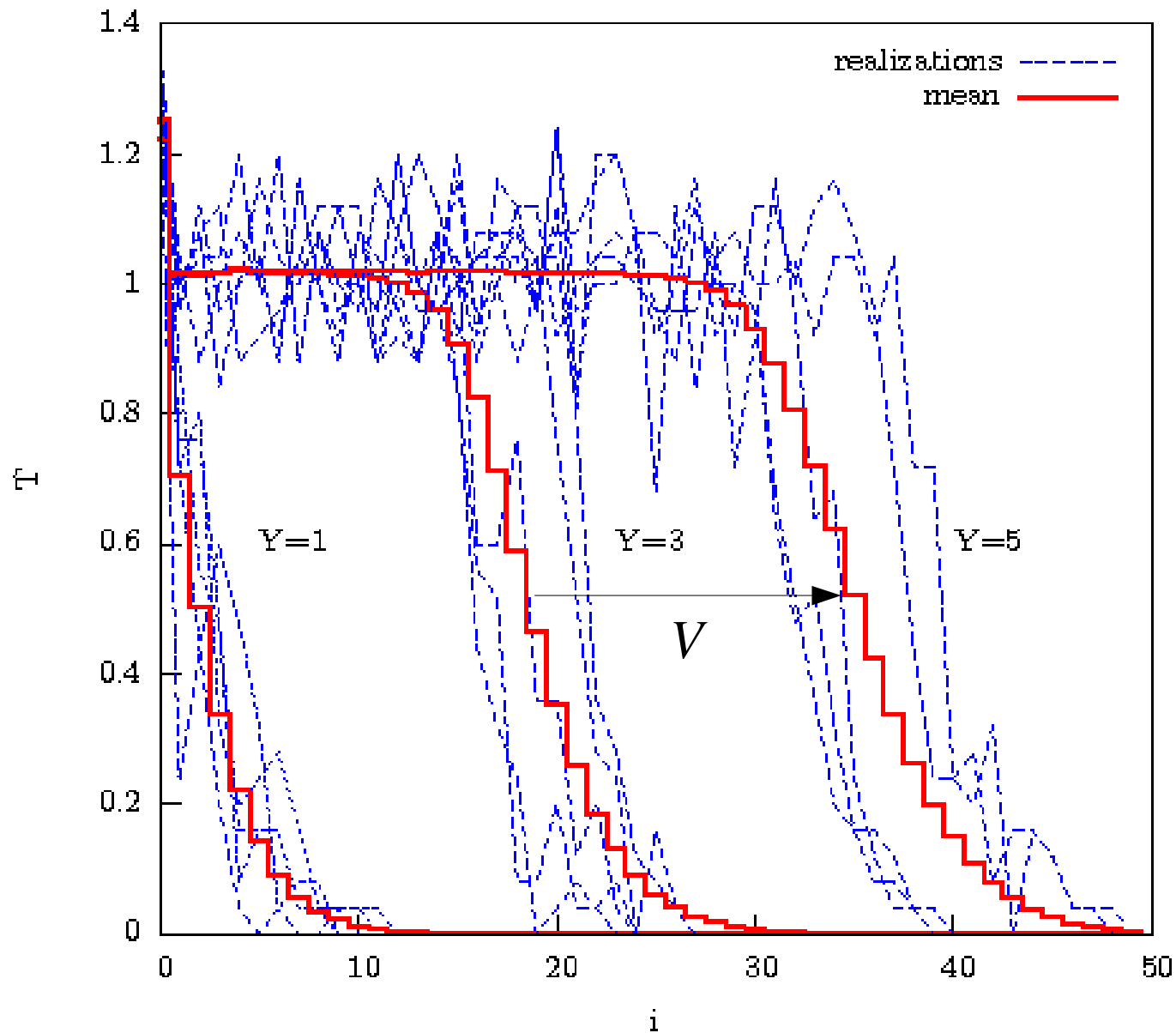
**We have 2 variables (r,b), and we keep the singularity structure of QCD**

# Traveling waves



One given impact parameter

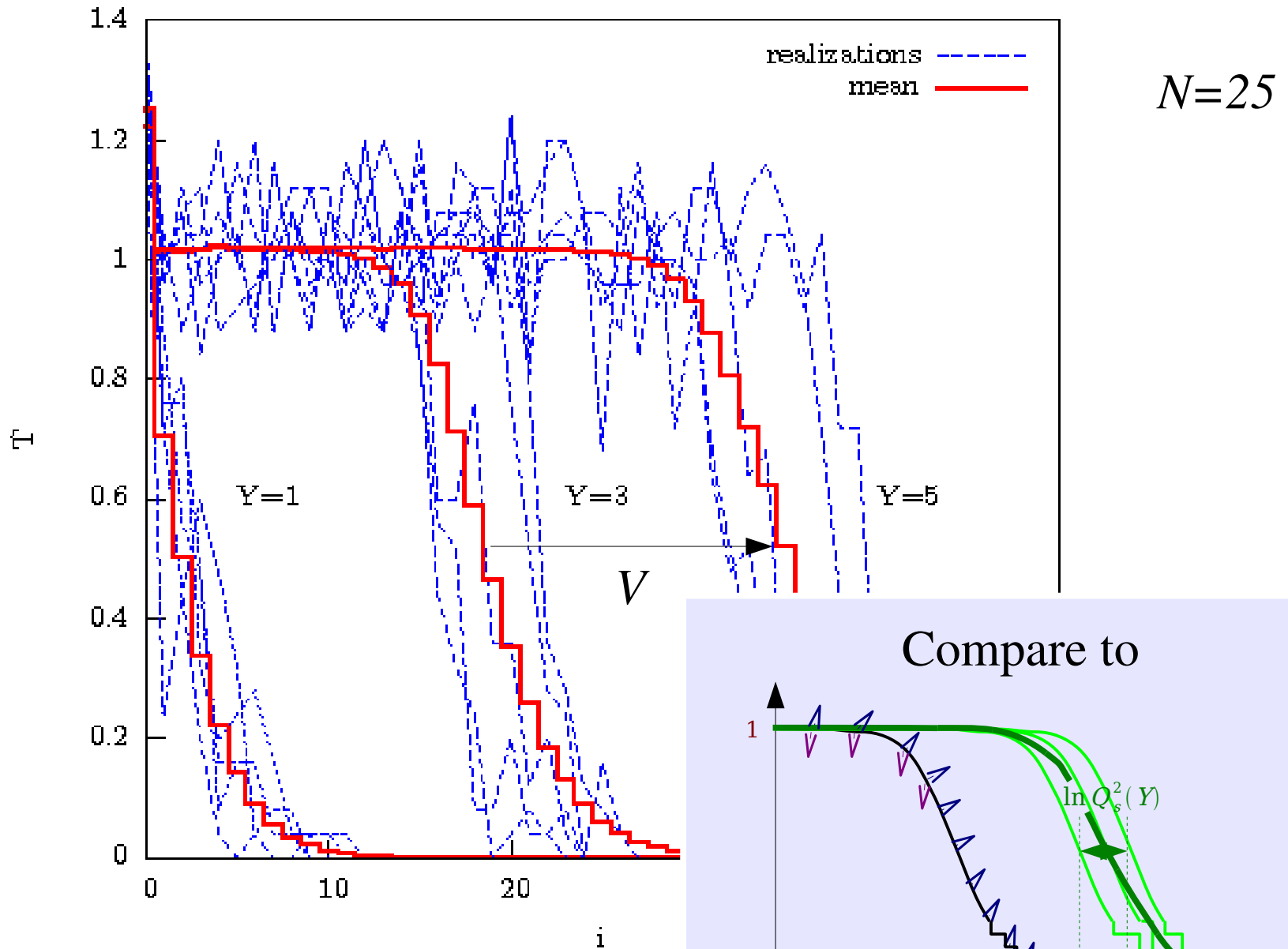
# Traveling waves



$N=25$

One given impact parameter

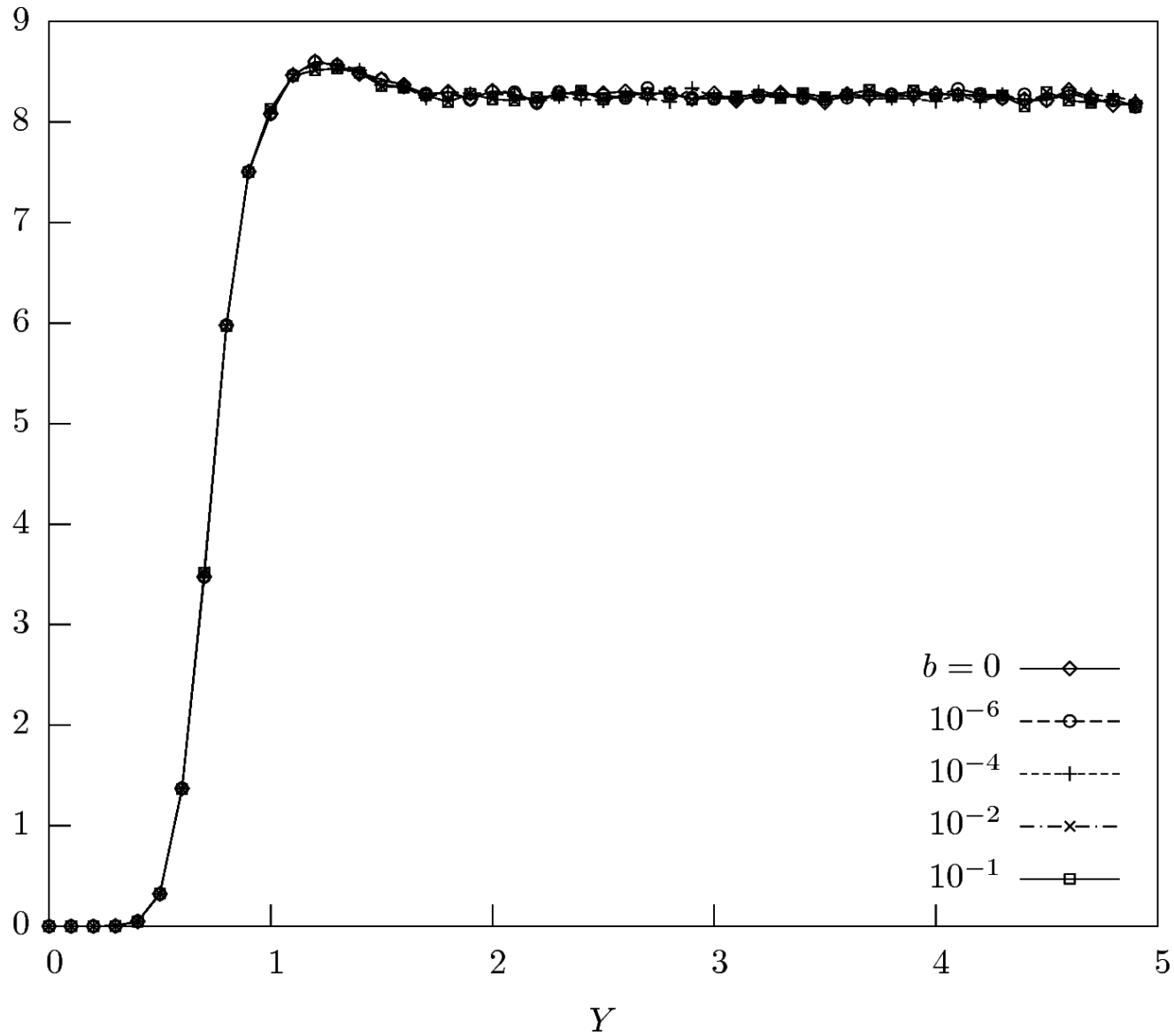
# Traveling waves



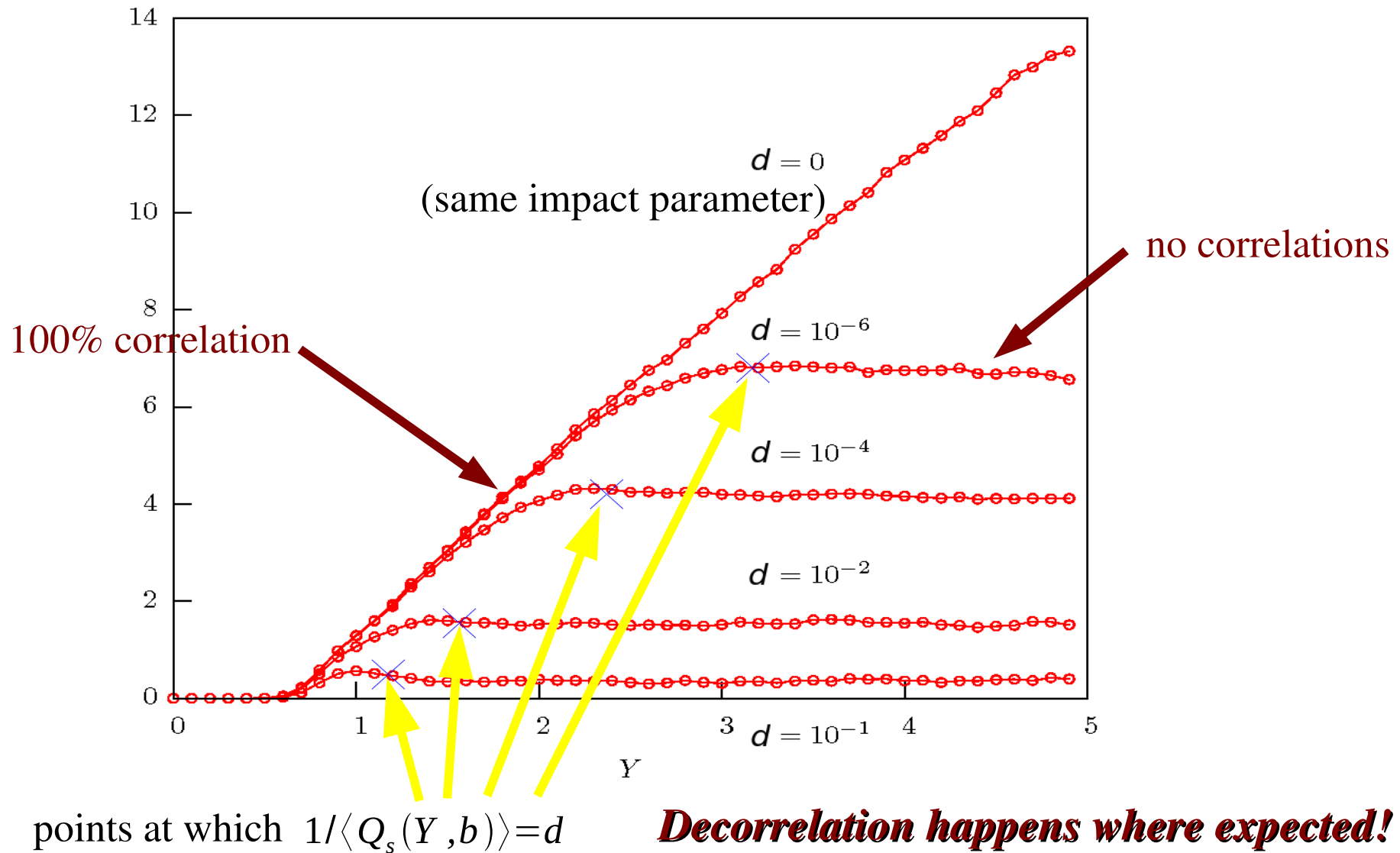
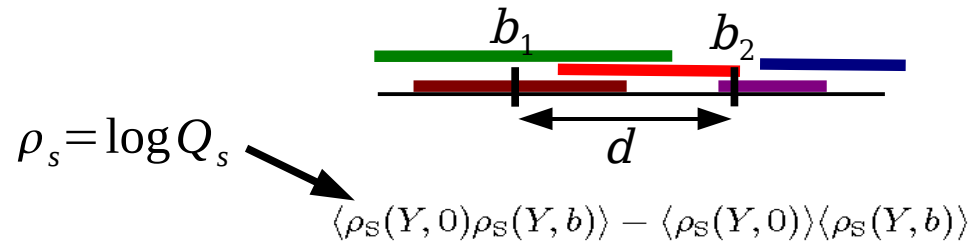
One given impact pa

# Traveling wave velocity

$$V = \langle \rho_S(Y + dY, b) - \rho_S(Y, b) \rangle / dY$$



# Correlation of the saturation scales



# *A more refined look*

*Comparison with a fixed impact-parameter version of the model*

The toy model is defined by its interval splitting rate

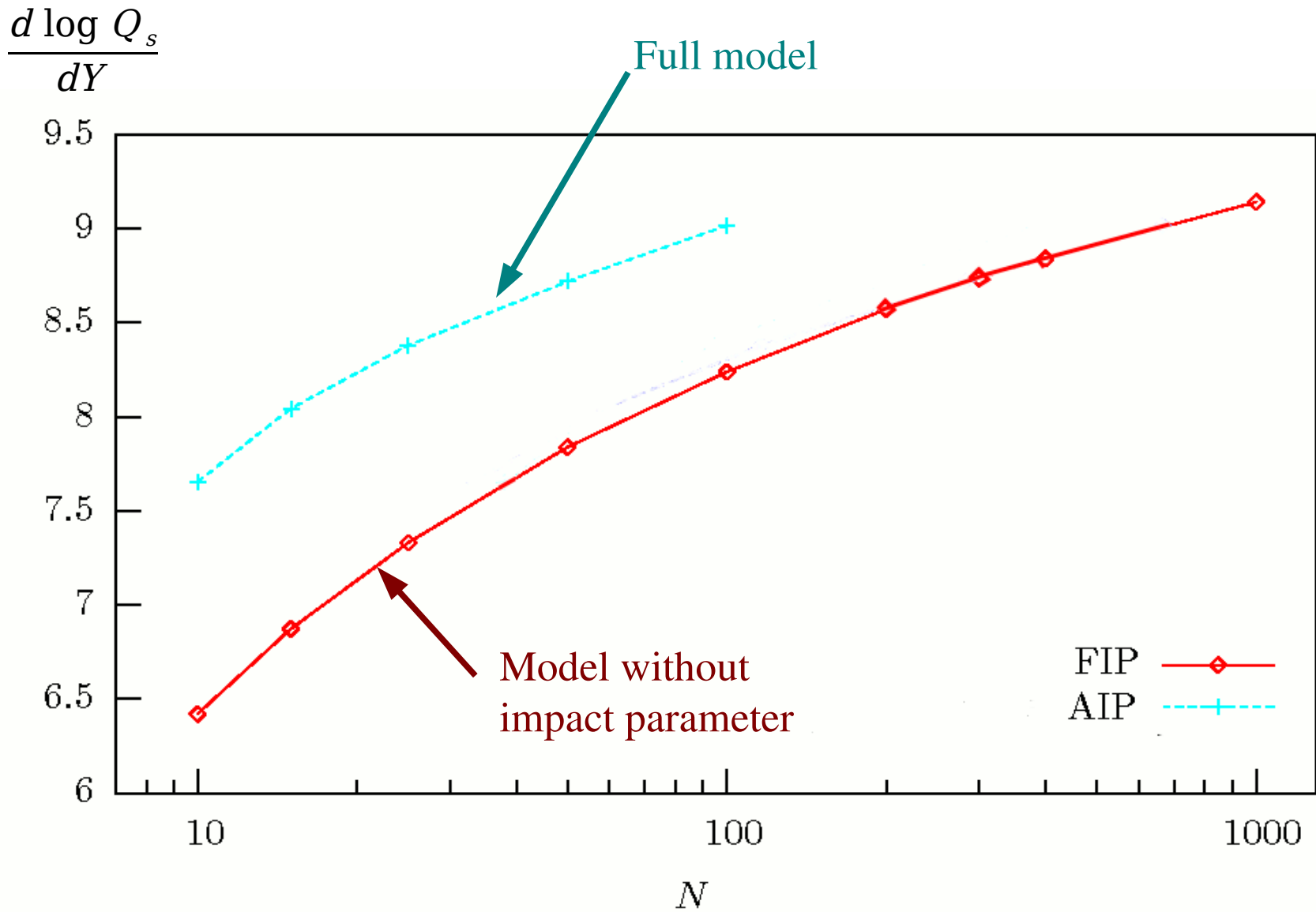
$$dP = \frac{|r_0|}{|r||r_0 - r|} dr dY \quad (+ \text{saturation condition})$$

which generates a distribution of sizes and impact parameters of intervals.

One may *discard the impact parameter dependence* (this implies a rescaling of the splitting rate) and get a true ***one-dimensional model*** for which only the size matters.

# *A more refined look*

*Comparison with a fixed impact-parameter version of the model*

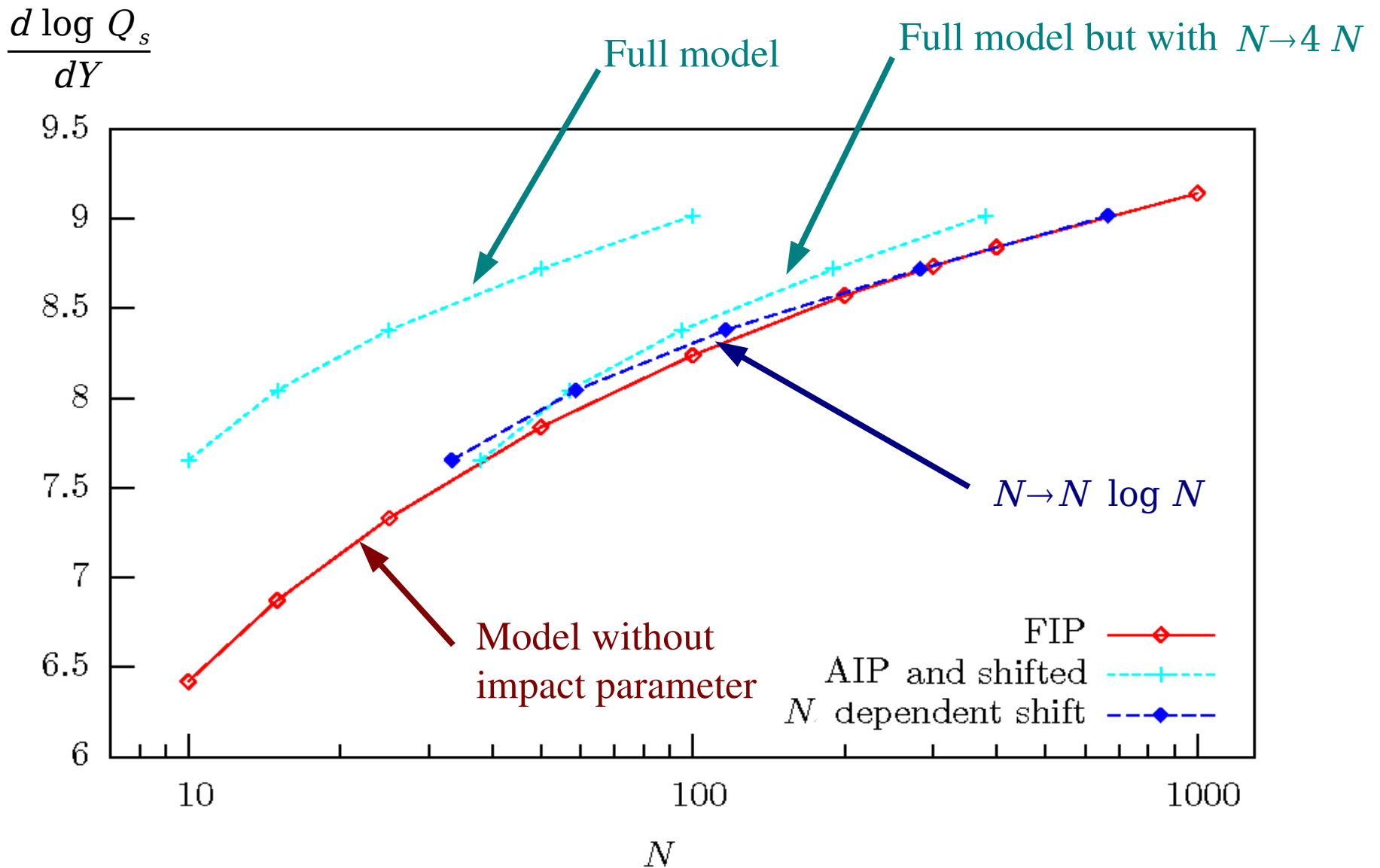


***Significant disagreement!***



# *A more refined look*

*Comparison with a fixed impact-parameter version of the model*



*The disagreement seems to amount to a mere rescaling of  $N$ !  
(=rescaling of the QCD coupling)*

# Summary

We have identified, from the physics, the universality class of high energy QCD as the one of *one-dimensional* reaction-diffusion processes, whose dynamics are governed by an equation of the form

$$\partial_{\bar{\alpha} Y} T = \chi(-\partial_{\ln k^2}) T - T^2 + \alpha_s \sqrt{T} \nu$$

We went back to the assumption that the QCD evolutions at different impact parameters decouple.

In a toy model, we have found that this is true.

However, a detailed comparison with a fixed-impact parameter version of the model shows some discrepancy, indicating that **the fixed-impact parameter model has more fluctuations.**