

Nucleon spin in perspective

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Epiphany 2009
Cracow, January 5 – 7, 2009

Introduction

- 1 Introduction
- 2 Longitudinal spin structure
- 3 Transversity
- 4 Gluon polarisation
- 5 Orbital angular momentum in the proton
- 6 Future
- 7 Outlook

A Beautiful Spin (after X. Ji)

- Born with troubles (Stern & Gerlach (1922) vs Goudsmit & Uhlenbeck (1925))
- Is due to space–time symmetry
- Fundamental concept
- Laboratory to explore physics beyond the SM, e.g.:
 - Muon “ $g - 2$ ” experiment @ BNL
 - Proton weak charge (Qweak exp @ JLAB)
 - Neutron EDM measurement ...

Tool to measure observables hard to obtain otherwise, e.g:

- Strangeness content of the nucleon from polarised parity–violating e–p scattering
- Electromagnetic form factors of the nucleon from the recoil polarisation
- Neutron density in large nuclei from parity–violating electron scattering
- and...

A Beautiful Spin (after X. Ji)...cont'd

- Probe to unravel the nonperturbative QCD dynamics, e.g.:
 - Nucleon spin-dependent structure functions, g_1 and g_2
 - Quark helicity ($\Delta q(x)$) and transversity ($\Delta_T q(x)$) distributions
 - Gluon polarisation, $\Delta g(x)$
 - Generalised Parton Distributions, GPD
 - Semi-Inclusive Deep Inelastic Scattering, SIDIS
 - (Generalised) Drell-Hearn-Gerasimov-... sum rule
 - Single spin asymmetries

Cross section asymmetries \implies structure functions

- In a full analogy to spin-averaged DIS where

$$\sigma \sim F_1(x) = \frac{1}{2} \sum_f e_f^2 q_f(x) \quad \text{and} \quad F_2(x) \approx 2xF_1$$

we now have: $\Delta\sigma = \sigma_{\rightarrow} - \sigma_{\leftarrow} \sim g_1(x) = \frac{1}{2} \sum_f e_f^2 \Delta q_f(x)$ and $g_2(x)$

where $\Delta q(x) = q^+(x) - q^-(x)$

- A direct observable, μ - p cross section asymmetry, $A^{\mu p}$:

$$A^{\mu d} = \frac{1}{fP_T P_B} \left(\frac{N_{\rightarrow} - N_{\leftarrow}}{N_{\rightarrow} + N_{\leftarrow}} \right); \quad f \sim 0.2, \quad P_T \sim 0.9, \quad P_B \sim -0.8$$

- is related to the longitudinal and transverse $\gamma^* p$ asymmetries:

$$\frac{A^{\mu p}}{D} = A_1^p + \eta A_2^p$$

- Longitudinal spin-dependent structure function:

$$g_1^p(x, Q^2) \approx A_1^p(x, Q^2) \frac{F_2^p(x, Q^2)}{2x(1 + R(x, Q^2))}$$

Cross section asymmetries \implies structure functions ...

- SIDIS asymmetry, e.g. A^- :

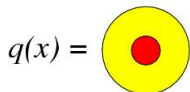
$$A^- = \frac{N_{\overleftarrow{\uparrow}}^{h^-} - N_{\overleftarrow{\downarrow}}^{h^-}}{N_{\overleftarrow{\uparrow}}^{h^-} + N_{\overleftarrow{\downarrow}}^{h^-}}$$

and

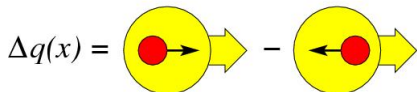
$$A_1^h = \frac{\sum_q e_q^2 (\Delta q D_q^h + \Delta \bar{q} D_{\bar{q}}^h)}{\sum_q e_q^2 (q D_q^h + \bar{q} D_{\bar{q}}^h)}$$

Partonic structure of the nucleon; distribution functions

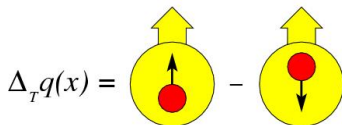
Three species of **twist-two** quark distributions in QCD (after integrating over the quark intrinsic k_t):



Quark momentum DF;
well known (unpolarised DIS $\rightarrow F_{1,2}(x)$).



Difference in DF of quarks with spin parallel or antiparallel to the nucleon's spin;
known, helicity (polarised DIS $\rightarrow g_1(x)$).



Difference in DF of quarks with spin parallel or antiparallel to the nucleon's spin in the transversely polarised nucleon;
unknown, transversity (polaris. SIDIS $\rightarrow A_{Collins}$)

In the nonrelativistic approach $\Delta_T q(x)$ identical with $\Delta q(x)$.

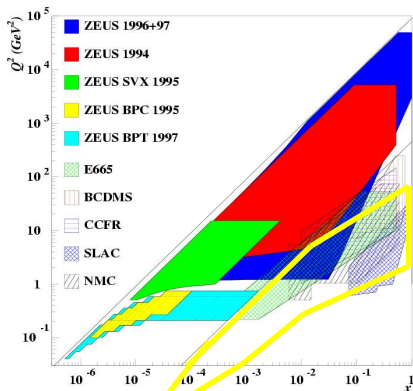
$\Delta_T q(x)$ are C-odd and chiral-odd; may only be measured as $\Delta_T q(x) \otimes \Delta_T D_q^q$.

If the k_t taken into account \implies 8 TMD appear; one, f_{1T}^\perp accessible through "Sivers asymmetry".

Recent and ongoing spin experiments

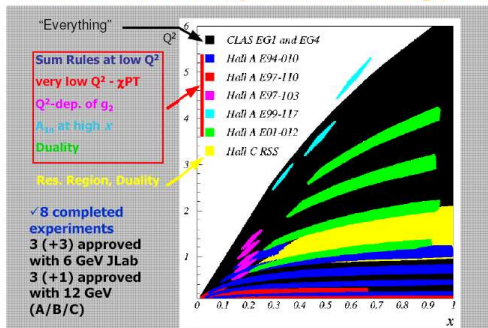
- CERN: EMC, SMC, COMPASS
- DESY: HERMES
- SLAC (completed): E142, E143, E154, E155, E156
- BNL (RHIC): STAR, PHENIX, BRAHMS

Acceptance of electroproduction experiments



COMPASS

Kinematics and Experimental Program



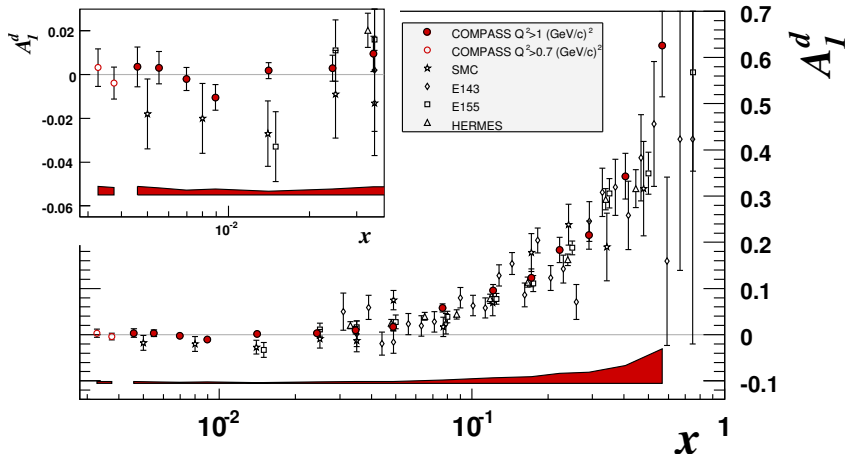
JLAB

Figure from: N. D'Hose, Villars 2004

Figure from: R. Carlini

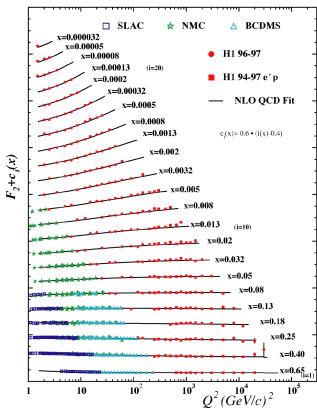
Longitudinal spin structure

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World data on $A_1^d(x)$ V.Yu. Alexakhin (COMPASS) *et al.* Phys. Lett. B **647** (2007) 8

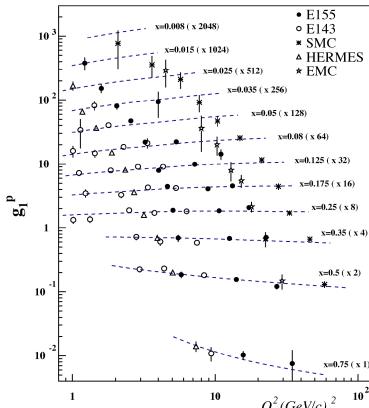
World data on the g_1^p (high energy experiments...)

World data on F_2^p



→ 50% of momentum
carried by gluons

World data on g_1^p



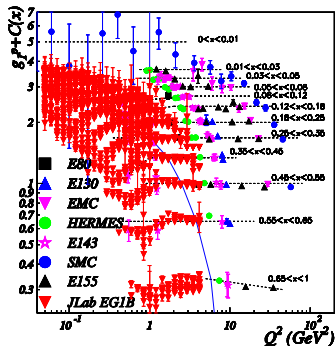
→ 20% of proton spin
carried by quark spin

Figure from R.Ent, DIS2006

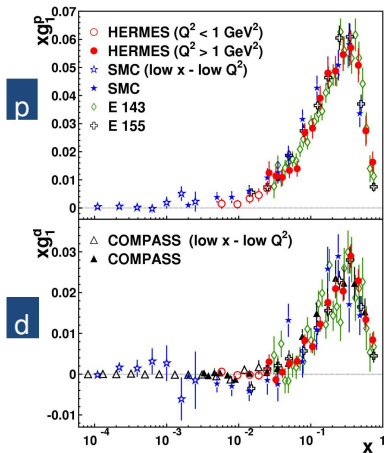
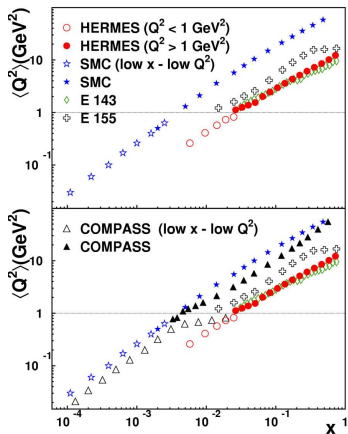
World data on the g_1^p (...and JLAB)World data on polarized structure
function $g_{1p}(x, Q^2)$

CLAS provides a large body of precise g_1 data in the DIS and transition regions that can be used to improve knowledge of twist-2 PDFs.

Phys. Rev.C75:035203, 2007
Phys. Lett. B 641, 11 (2006)



World data on the nucleon g_1

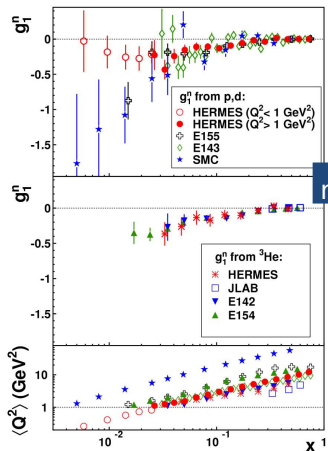


World data on the nucleon g_1 ...cont'd

from p and d

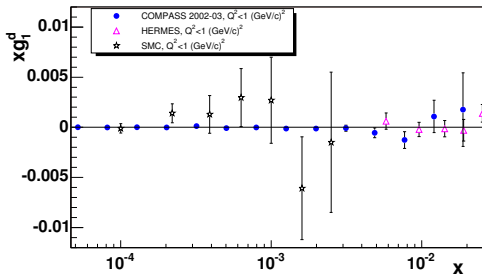
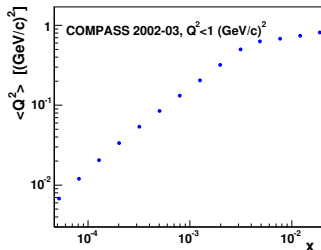
$$g_1^n = \frac{2}{1 - \frac{3}{2}\omega_D} g_1^d - g_1^p$$

from ^3He



g_1^d in the nonperturbative ($Q^2 < 1 \text{ (GeV/c)}^2$) region)

V.Yu. Alexakhin (COMPASS) *et al.* Phys. Lett. B **647** (2007) 330

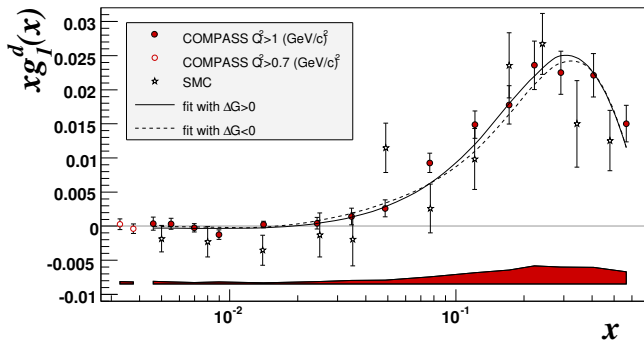


- Order of magnitude improvement over the statistical precision of the SMC.
- Interplay between perturbative and nonperturbative mechanisms.
- Spin effects in g_1^d at low x and Q^2 absent ?

COMPASS QCD analysis of inclusive g_1^d

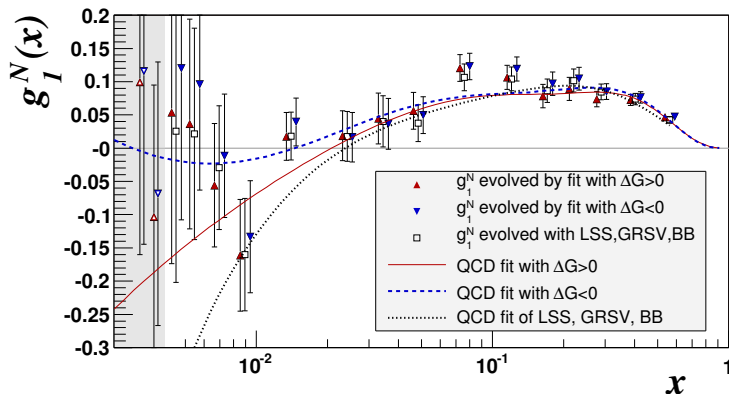
V.Yu. Alexakhin (COMPASS) *et al.* Phys Lett B647 (2007) 8

- Two programs: DGLAP evolution of structure functions and evolutions of moments
- NLO \overline{MS} scheme
- World data: 9 experiments, 230 data points (43 from COMPASS)
- Two solutions, $\Delta G > 0$ and $\Delta G < 0$ describe data equally well.



Quark polarisation from COMPASS data only (@ $Q^2 = 3 \text{ GeV}^2$):

$a_0 = 0.35 \pm 0.03(\text{stat.}) \pm 0.05(\text{syst.})$ and gluon polarisation: $|\Delta G| \approx 0.2 - 0.3$

COMPASS QCD analysis of g_1^d ...cont'dV.Yu. Alexakhin (COMPASS) *et al.* Phys Lett B647 (2007) 8

COMPASS g_1^N evolved to $Q^2 = 3 \text{ GeV}^2$; LSS, GRSV, BB are NLO fits to world (but no COMPASS) data.

Low x data prefer $\Delta G < 0$??? Sign of ΔG not fixed by the g_1 measurements...

Gluon polarisation from QCD evolution

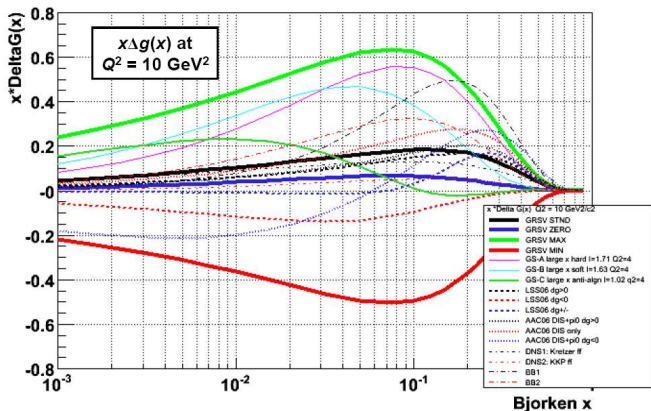


Figure from C. Gagliardi, DIS2008

Carl Gagliardi – DIS2008 – Jets in pp at RHIC

NLO QCD analysis of world data...cont'd

TABLE II: First moments $\Delta f_j^{1,[x_{\min}^{-1}]}$ at $Q^2 = 10 \text{ GeV}^2$.

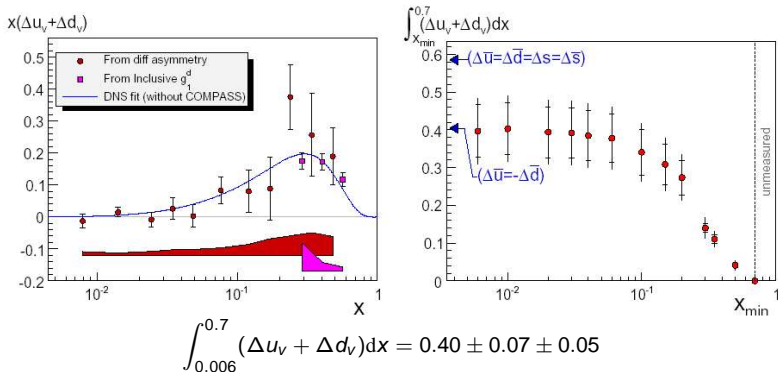
	$x_{\min} = 0$	$x_{\min} = 0.001$	
	best fit	$\Delta\chi^2 = 1$	$\Delta\chi^2/\chi^2 = 2\%$
$\Delta u + \Delta\bar{u}$	0.813	0.793 $^{+0.011}_{-0.012}$	0.793 $^{+0.028}_{-0.034}$
$\Delta d + \Delta\bar{d}$	-0.458	-0.416 $^{+0.011}_{-0.009}$	-0.416 $^{+0.035}_{-0.025}$
$\Delta\bar{u}$	0.036	0.028 $^{+0.021}_{-0.020}$	0.028 $^{+0.059}_{-0.059}$
$\Delta\bar{d}$	-0.115	-0.089 $^{+0.029}_{-0.029}$	-0.089 $^{+0.090}_{-0.080}$
$\Delta\bar{s}$	-0.057	-0.006 $^{+0.010}_{-0.012}$	-0.006 $^{+0.028}_{-0.031}$
Δg	-0.084	0.013 $^{+0.106}_{-0.120}$	0.013 $^{+0.702}_{-0.314}$
$\Delta\Sigma$	0.242	0.366 $^{+0.015}_{-0.018}$	0.366 $^{+0.042}_{-0.062}$

- First moment of the singlet distribution ~ 0.25 !
- Gluon polarisation small (with large errors!)

Sea quark polarisation

M. Alekseev *et al.* (COMPASS), Phys. Lett. **B660** (2008) 458.

- Difference asymmetry: $A^{h^+-h^-} : A_d^{\pi^+-\pi^-} = A_d^{K^+-K^-}$
- At LO, the fragmentation functions drop out



- Unmeasured regions contribute negligibly.
- Non-symmetric sea preferred ?
- Next step: determine Δs from K^\pm asymmetries.

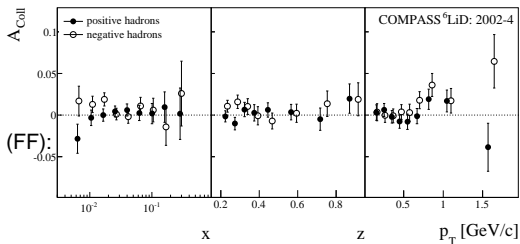
Transversity

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Results for the Collins and Sivers asymmetries

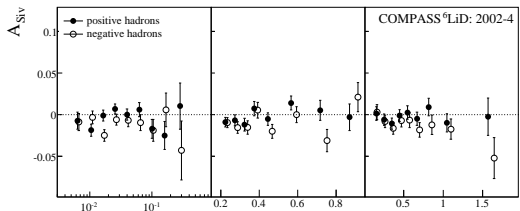
Deuteron target; all hadrons: positive and negative

E.S. Ageev *et al.* (COMPASS) Nucl.Phys.B **765**(2007) 31



Collins asymmetries very small.
These data + Hermes + Belle

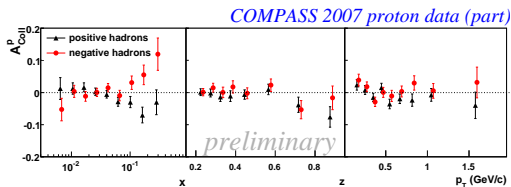
$$\Rightarrow \Delta_T u + \Delta_T d \sim 0$$



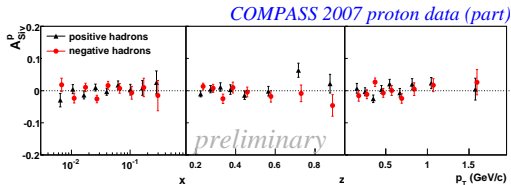
Sivers asymmetries very small.
S.Brodsky & S. Gardner (2006):
no gluon orbital angular
momentum in the nucleon?

Results for the Collins and Sivers asymmetries...cont'd

Proton target; all hadrons: positive and negative



Asymmetries nonzero at $x \gtrsim 0.1$.



Sivers asymmetries very small.

Gluon polarisation

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Nucleon spin decomposition

$$\frac{\hbar}{2} = J_q + J_g = \left(\frac{1}{2} \Delta \Sigma + L_q \right) + \left(\Delta G + L_g \right)$$

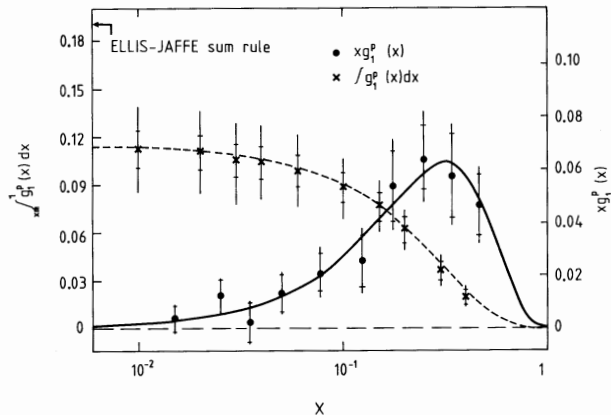
- Observe! each term is μ^2 dependent;
decomposition $J_g = \Delta G + L_g$ is NOT gauge-invariant.
- EMC (1988): $a_0 = \Delta \Sigma = 0.12 \pm 0.09 \pm 0.14$ (expected: ~ 0.6 if $\Delta s = 0$).
Here $\Delta \Sigma = \Delta u + \Delta d + \Delta s$ and $\Delta q = \int \Delta q(x) dx$, $\Delta G = \int \Delta G(x) dx$
- COMPASS @ 3 GeV²: $a_0 = 0.35 \pm 0.03 \pm 0.05$
- But as a consequence of the “axial anomaly” (axial vector current not conserved) the measured quantity is:

$$a_0(Q^2) = \Delta \Sigma^{AB} - \left(\frac{n_f \alpha_s}{2\pi} \right) \Delta G(Q^2)$$

and the “spin crisis” can be solved ($\Delta \Sigma \sim 0.6$) if $\Delta G \sim 2.2$ (and $L \sim -2$) at $Q^2 = 3 \text{ GeV}^2$.

- Impressive spin-off since 1988: SLAC (E142, E143, E155, E156), SMC, HERMES, JLAB, COMPASS, RHIC Spin.
- **Need to measure ΔG (and L)!**

“Spin puzzle”: 20 years



European Muon Collaboration, J.Ashman *et al.* Phys. Lett. **B206** (1988) 364

$$\Gamma_1^p = 0.123 \pm 0.013 \pm 0.019$$

$$\Delta\Sigma = a_0 = 0.12 \pm 0.17$$

Ways of measuring the ΔG

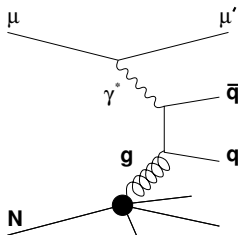
- Scaling violation of g_1 (QCD fits, world data).
- Direct measurements
 - Cross section asymmetry for the **photon–gluon fusion (PGF)** with subsequent fragmentation into the **charm mesons** (max. @ low Q^2 , perturbative scale: e.g. mass of the charm quark).
 - Cross section asymmetry for the **photon–gluon fusion (PGF)** with subsequent fragmentation into a **pair of hadrons of large p_T** , separately for low- and high Q^2 (perturbative scale: e.g. p_T).

Direct $\Delta G/G$ measurements

Mechanism employed: photon–gluon fusion. **Observable:** asymmetry in the hadron production

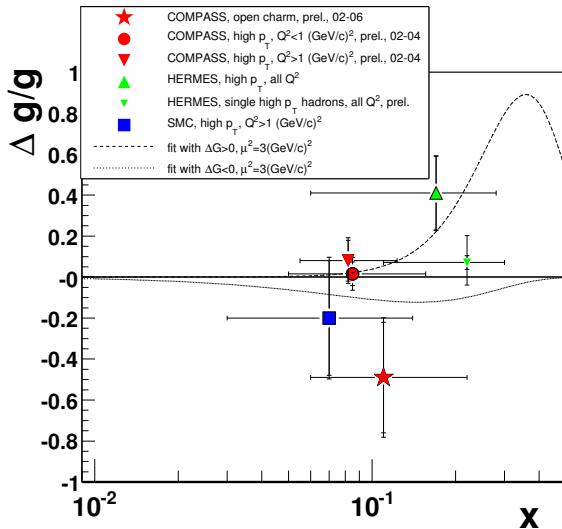
$$A_{\gamma N}^{PGF} = \frac{\int d\hat{s} \Delta\sigma^{PGF} \Delta G(x_g, \hat{s})}{\int d\hat{s} \sigma^{PGF} G(x_g, \hat{s})}$$

$$\approx \langle a_{LL}^{PGF} \rangle \frac{\Delta G}{G}$$



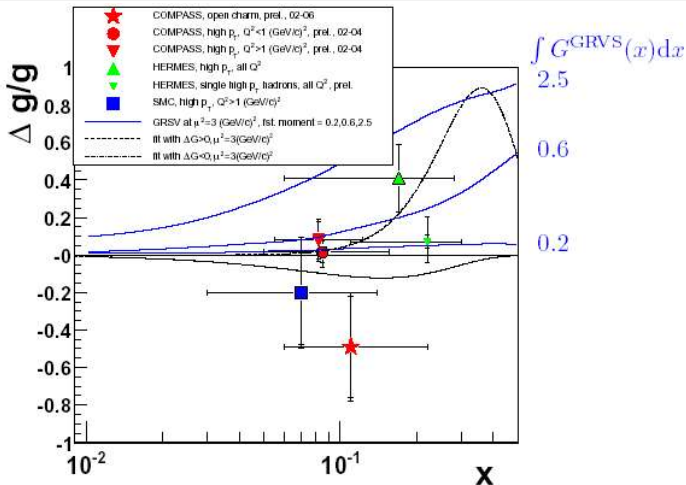
- If $q \equiv c \implies$ a pair of charmed mesons (we demand only one) in the final state;
 - measurement difficult (low statistics),
 - NLO corrections may be important.
- If $q \equiv u, d, s \implies$ a pair of jets or (in COMPASS) of high- p_T hadrons;
 - measurement simple (high statistics),
 - strong physics background, MC – dependence.

Summary of the gluon polarisation measurements



At $x_g \sim 0.1$, $\Delta G/G$ is compatible with zero! Qualitative agreement with RHIC results.

Summary of the gluon polarisation measurements



- Restoration of $\Delta\Sigma=0.6$ via the axial anomaly improbable.
- Independent measurement of orbital momentum necessary!

Orbital angular momentum in the proton

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Ways of constraining the orbital angular momentum

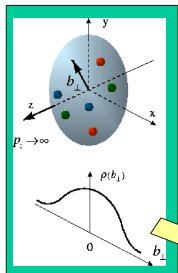
- Extract from the full, 3D description of the proton \implies GPDs
- Measurement of asymmetries sensitive to e.g. the k_T of the partons (Sievers, Cahn ???)
- QCD calculations on the lattice
- X. Ji: evolution equations for J_q and J_g give the asymptotic solutions:

$$J_q(\mu^2 \rightarrow \infty) = \frac{1}{2} \frac{3n_f}{16 + 3n_f}, \quad J_g(\mu^2 \rightarrow \infty) = \frac{1}{2} \frac{16}{16 + 3n_f}$$

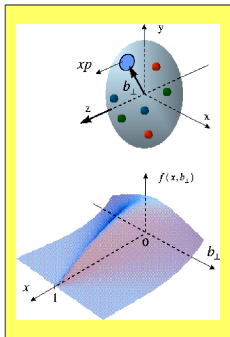
L cannot be negligible!

3D picturing of the proton *via* GPD

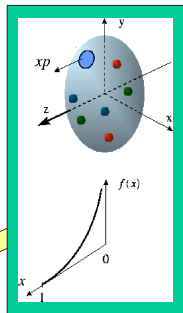
D. Mueller, X. Ji, A. Radyushkin, A. Belitsky, ...
 M. Burkardt, ... Interpretation in impact parameter space



Proton form factors,
transverse charge &
 current densities

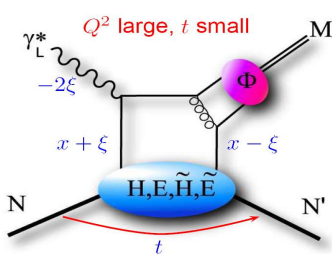


Correlated quark momentum
 and helicity distributions in
 transverse space - **GPDs**



Structure functions,
 quark **longitudinal**
 momentum & helicity
 distributions

Access GPD through the DVCS mechanism



- Four GDPs ($H, E, \tilde{H}, \tilde{E}$) for each flavour and for gluons
- Factorisation proven for σ_L only
- All depend on 3 variables: x, ξ, t ; DIS @ $\xi = t = 0$
- H, \tilde{H} conserve nucleon helicity
 E, \tilde{E} flip nucleon helicity
- H, E refer to unpolarised distributions
- $H^q(x, 0, 0) = q(x)$, $\tilde{H}^q(x, 0, 0) = \Delta q(x)$
 \tilde{H}, \tilde{E} refer to polarised distributions

Figure from A. Sandacz, EINN, 2007

- H, E accessed in vector meson production via A_{UT} asymmetries
- \tilde{H}, \tilde{E} accessed in pseudoscalar meson production via A_{UT} asymmetries
- All 4 accessed in DVCS (γ production) in $A_C, A_{LU}, A_{UT}, A_{UL}$
- Integrals of $H, E, \tilde{H}, \tilde{E}$ over x give Dirac-, Pauli-, axial vector- and pseudoscalar vector form factors resp.
- **Important:** $J_z^q = \frac{1}{2} \int dx x [H^q(x, \xi, t=0) + E^q(x, \xi, t=0)] = \frac{1}{2} \Delta \Sigma + L_z^q$ (X. Ji)

DVCS data taken or expected soon; future projects

- Data being analysed:
 - Z1 and ZEUS: cross section for $\rho^0, \phi, J/\psi, \Upsilon$ production
 - HERMES: cross section for $\rho^0, \phi, \text{BCA, BSA, TSA } (\pi^+)$
 - CLAS: cross section for $\rho^+, \phi, \pi^+, \pi^0, \eta$
- Data taken or expected
 - DVCS Coll. (JLAB, Hall A): cross section for π^0 (2009)
 - CLAS: BSA, TSA (2009)
 - COMPASS: muoproduction of ρ^0 on the proton (2007 data)
- Future
 - COMPASS upgrade for DVCS, DVMP (\sim 2010)
 - JLAB 12 GeV upgrade (\gtrsim 2014)
 - PANDA @ FAIR (\gtrsim 2014)
 - EIC/ELIC (\gtrsim ??)
- First results constraining E and H (and thus orbital momentum) and DVMP models; L_q close to zero ?

Orbital angular momentum in the proton: results from lattice QCD

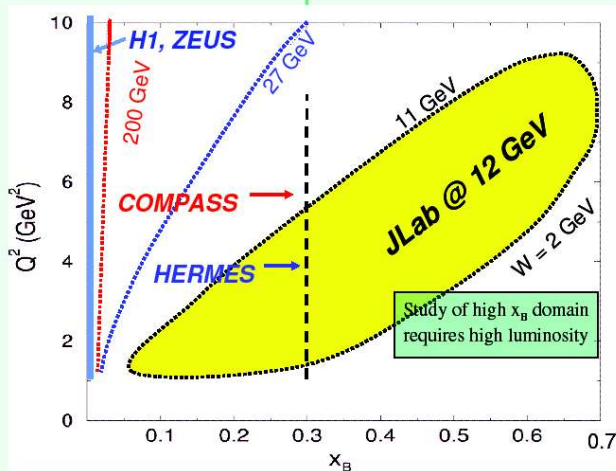
- In the sum rule: $\frac{1}{2} = J_q + J_g = \frac{1}{2}\Delta\Sigma + L_q + \Delta G + L_g$
- total spin: $J_q^u = 0.22 \pm 0.02$, $J_q^d = 0.00 \pm 0.02$
- OAM : $L^u = 0.20 \pm 0.04$, $L^d = -0.20 \pm 0.04$
 $\implies L_q = 0 ?$
- Errors do not contain the systematic part (hard to estimate) !
- Results on L_q in line with the DVCS results of HERMES, JLAB.

Future

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CEBAF 12 GeV upgrade

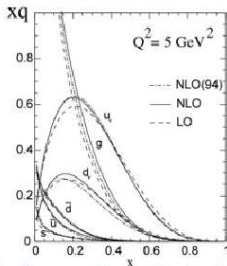
*Deeply Virtual Exclusive Processes -
Kinematics Coverage of the 12 GeV Upgrade*



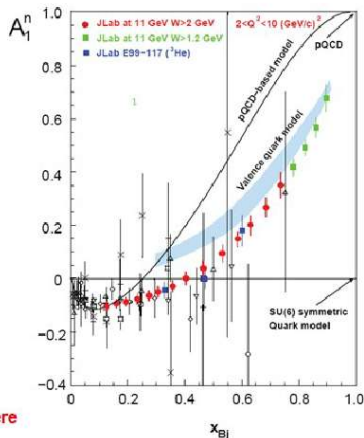
Future: high x structure functions at JLAB 12 GeV

REQUIRES:

- High beam polarization
- High electron current
- High target polarization
- Large solid angle spectrometers



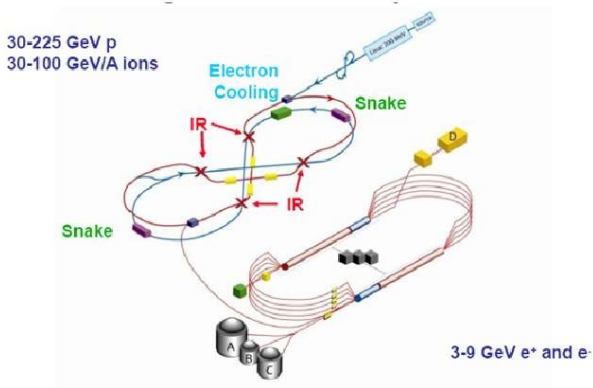
12 GeV will access the regime ($x > 0.3$), where valence quarks dominate



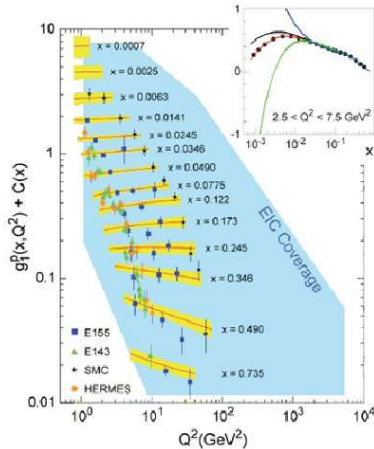
Slide from H. Montgomery, SPIN 2008

Future: Electron-Ion Collider, EIC

- LHeC @ CERN
- eRHIC @ BNL
- ELIC @ JLAB
- MANUEL @ FAIR (GSI)



Future: Electron–Ion Collider, EIC,...cont'd

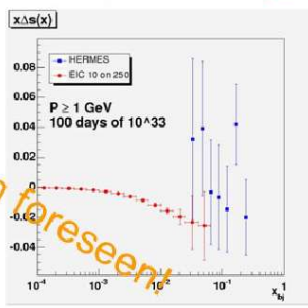
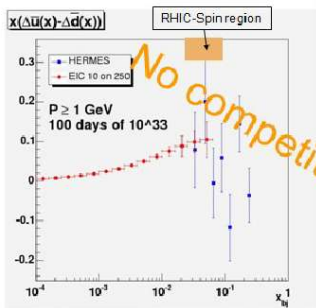
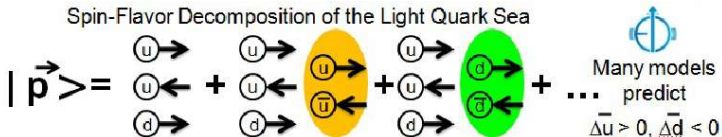


Slide from H. Montgomery, SPIN 2008



Future: Electron-Ion Collider, EIC,...cont'd

Spin-Flavor Decomposition of the Light Quark Sea



No competition foreseen!

Slide from H. Montgomery, SPIN 2008



Outline

- 1 Introduction
- 2 Longitudinal spin structure
- 3 Transversity
- 4 Gluon polarisation
- 5 Orbital angular momentum in the proton
- 6 Future
- 7 Outlook**

Outlook

- We have learned a lot about the polarised PDFs (helicity, transversity) in the last 20 years.
- “Spin puzzle”: restoration of $\Delta\Sigma = 0.6$ via the axial anomaly improbable. Significant orbital angular momentum in the proton expected; must find a way to expose it \implies DVCS!
- Much experimental and theoretical progress in analysing the physics mechanism of the Single-Spin-Asymmetries.
- More data await analysis (HERMES, COMPASS, RHIC, JLAB); new data come (COMPASS, RHIC); new experiments soon (COMPASS II (2010), JLAB 12 (2014)).
- **Polarised EIC badly needed!**

THANKS, Jan!



SPARE

Spin applications...

1

Fluctuations and Market Friction in Financial Trading

Bernd Rosenow

Institut für Theoretische Physik, Universität zu Köln, D-50937 Köln, Germany

(May 15, 2006)

We study the relation between stock price changes and the difference in the number of sell and buy orders. Using a soft spin model, we describe the price impact of order imbalances and find an analogy to the fluctuation-dissipation theorem in physical systems. We empirically investigate fluctuations and market friction for a major US stock and find support for our model calculations.

PACS numbers: 05.45.Tp, 89.90.+n, 05.40.-a, 05.40.Fb

nat/0107018 v2 6 Jul 2001

The unpredictable up and down movements in the stock market have always captured the interest and imagination of investors. The scientific investigation of these phenomena started with Bachelier's comparison of stock price dynamics with a random walk [1]. This study has been refined in many respects [2]. Interest has been devoted to the precise shape of the distribution of returns (difference of the logarithm of stock prices at different times), which is characterized by a high probability for large fluctuations [3-6].

After it was realized that large price fluctuations tend to cluster together in time, stock price returns were described by models of volatility (standard deviation of returns) changing in time [7,8]. This effect is captured in time series models, in which the volatility at a given time depends on the magnitude of previous returns [9,10]. It has been actively investigated how the stochastic properties of price dynamics can be related to the market microstructure, i.e. the rules and motivations according

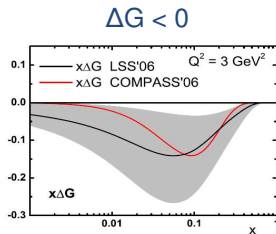
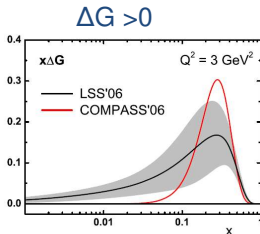
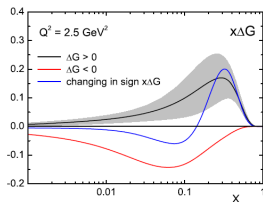
to sell orders (order imbalance), which acts as an external force. We study the dependence of stock price changes on order imbalance empirically by using the method of data analysis and some of the results of [21,22]. We find that the empirical results agree well with our model.

Model calculation: The observable quantity we are interested in is the logarithmic stock price changes within a time interval Δt

$$G_{\Delta t}(t) = \ln S(t) - \ln S(t - \Delta t), \quad (1)$$

where $S(t)$ is the price of a given stock at time t . Transaction prices at a stock exchange lie usually in a finite interval between the bid price (the price a trader offers to pay for a stock) and the ask price (the price at which a dealer is willing to buy the stock). In addition, the historical prices studied take only discrete (tick) values. This motivates to model price changes by a spin model, and for the virtue of easier analytical calculations by a

Gluon polarisation from QCD evolution, cont.'d



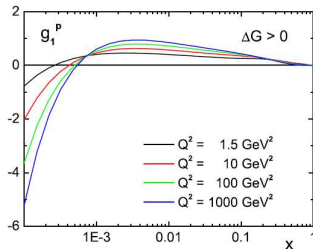
Can we ever tell the sign of ΔG ?

Figures from E. Leader, DIS2008

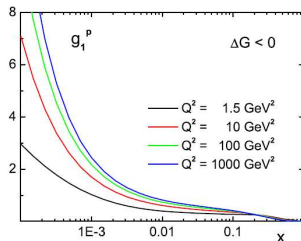
Gluon polarisation from QCD evolution, cont.'d

... except at an ep collider ?

LSS06, $\Delta G > 0$



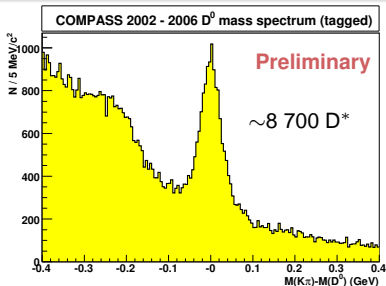
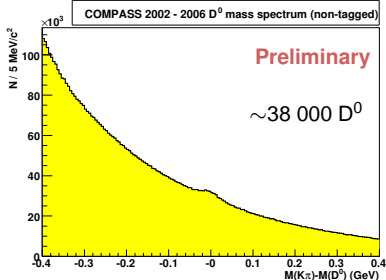
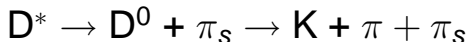
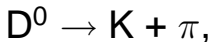
LSS06, $\Delta G < 0$





Figures from E. Leader, DIS2008

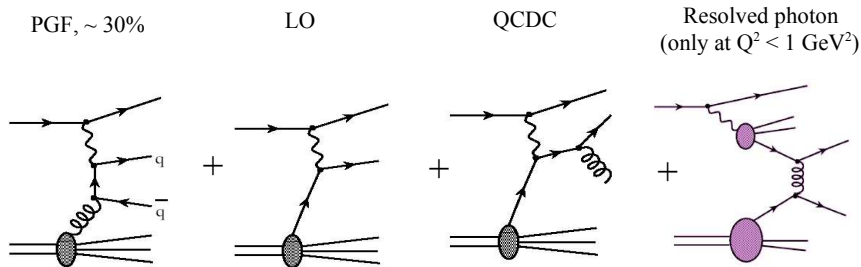


Direct $\Delta G/G$ measurements; open charm production



- Choose $D^0 \rightarrow K\pi$ ($BR \sim 4\%$); pions and kaons identified by RICH.
- Clean sample of the PGF events (but low statistics); little physics background.
- Combinatorial background significantly reduced for the $D^* \rightarrow D^0 + \pi_S \rightarrow K + \pi + \pi_S$.
- Charm in the nucleon neglected.
- Weak dependence on the MC in the analysis.
- **A weighting method used to optimise the $\Delta G/G$ extraction**  **C. Quintans' talk** 

Direct $\Delta G/G$ measurements; high p_T hadrons



$$A_{LL}^{2h}(x) = R_{pgf} \cdot a_{LL}^{pgf} \cdot \frac{\Delta G}{G}(x_g) + R_{LO} \cdot D \cdot A_1^{LO}(x) + R_{QCDQ} \cdot a_{LL}^{QCDQ} \cdot A_1^{LO}(x_c)$$

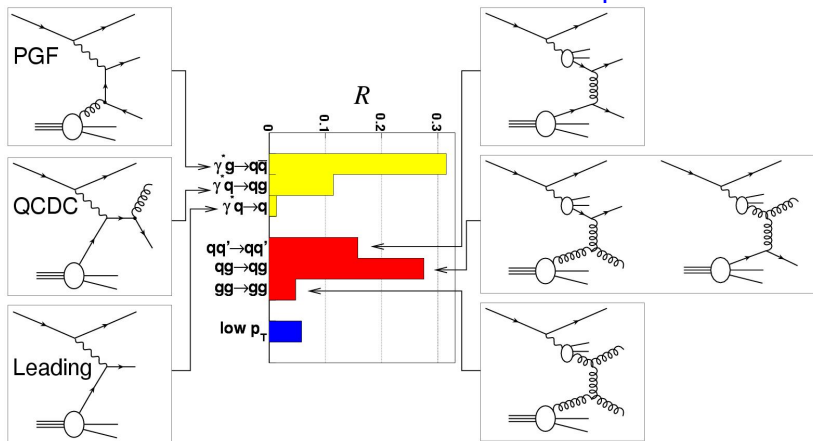
$\frac{\Delta G}{G}$ evaluated from $A_{LL}^{2h} = A_{meas}/(P_b P_T f)$ and from inclusive asymmetries, $A_1 = A_{meas}^{incl}/D$ if **contributions** of PGF and of background processes (LO, QCDQ) **taken from MC**

⇒ C. Quintans' talk.

Direct $\Delta G/G$ measurements; high p_T hadrons @ $Q^2 < 1 \text{ GeV}^2$

E.S. Ageev (COMPASS) *et al.* Phys. Lett. B **633** (2006) 25

Resolved photons



Summary of the gluon polarisation measurements

...cont'd

$$\frac{1}{2} = \frac{1}{2}\Delta\Sigma + \Delta G + L$$

Are we approaching the solution of the “proton spin puzzle”?

- Restoration of $\Delta\Sigma=0.6$ via the axial anomaly improbable.
- Global, consistent NLO analysis of ΔG needed.
- Independent measurement of L necessary.

Properties of transversity

Properties of $\Delta_T q(x)$:

- is chiral-odd \implies hadron(s) in final state needed to be observed
- simple QCD evolution since no gluons involved
- related to GPD
- sum rule for transverse spin
- first moment gives “tensor charge” (now being studied on the lattice)

Asymmetry measured e.g. via the Collin’s asymmetry (asymmetry in the distribution of hadrons):

$$N_h^\pm(\phi_c) = N_h^0 [1 \pm p_T D_{NN} A_{Coll} \sin \phi_c]$$

which in turn gives at LO:

$$A_{Coll} \sim \frac{\sum_q e_q^2 \cdot \Delta_T q \cdot \Delta_T^0 D_q^h}{\sum_q e_q^2 \cdot q \cdot D_q^h}$$

But **transverse fragmentation functions $\Delta_T^0 D_q^h$** needed to extract $\Delta_T q(x)$ from the Collin’s asymmetry! Recently those FF measured by BELLE.

Properties of the Sivers process:

it is related to L_q in the proton. **Fundamental !**

Link to DIS and elastic form factors

DIS at $\xi=t=0$
 $H^q(x,0,0) = q(x)$
 $\tilde{H}^q(x,0,0) = \Delta q(x)$

Form factors (sum rules)

$$\int_{-1}^1 dx \sum_q [H^q(x, \xi, t)] = F_1(t) \text{ Dirac f.f.}$$

$$\int_{-1}^1 dx \sum_q [E^q(x, \xi, t)] = F_2(t) \text{ Pauli f.f.}$$

$$\int_{-1}^1 dx \tilde{H}^q(x, \xi, t) = G_{A,q}(t), \quad \int_{-1}^1 dx \tilde{E}^q(x, \xi, t) = G_{P,q}(t)$$

$$H^q, E^q, \tilde{H}^q, \tilde{E}^q(x, \xi, t)$$

Angular Momentum Sum Rule

$$J^q = \frac{1}{2} - J^G = \frac{1}{2} \int_{-1}^1 x dx [H^q(x, \xi, 0) + E^q(x, \xi, 0)]$$

X. Ji, Phys.Rev.Lett.78,610(1997)

Deeply Virtual Compton Scattering

$$ep \rightarrow e\gamma$$

Kinematics

