

Compact galactic sources of cosmic radiation

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Summary

1. Pulsar emission
2. Crab nebula
3. Pulsar wind model
4. Neutron stars with ultra-strong magnetic fields

1. Pulsar emission

a) Pulsar energy output

Standard radio-pulsars

- rotating neutron stars, periods $P \sim 10ms - 1s$,
- mass $M \sim 1M_{\odot}$,
- spin-down rates $\dot{P} \sim 10^{-11} - 10^{-14} s s^{-1}$,
- magnetic field $B \sim 10^{12} G$.

Spin-down luminosity

$$\dot{E} \equiv \dot{E}_{rot} = I\omega\dot{\omega} \sim 10^{35} - 10^{39} ergs^{-1}$$

- the moment of inertia, $I \sim 10^{45} gcm^2$, from neutron

star models,

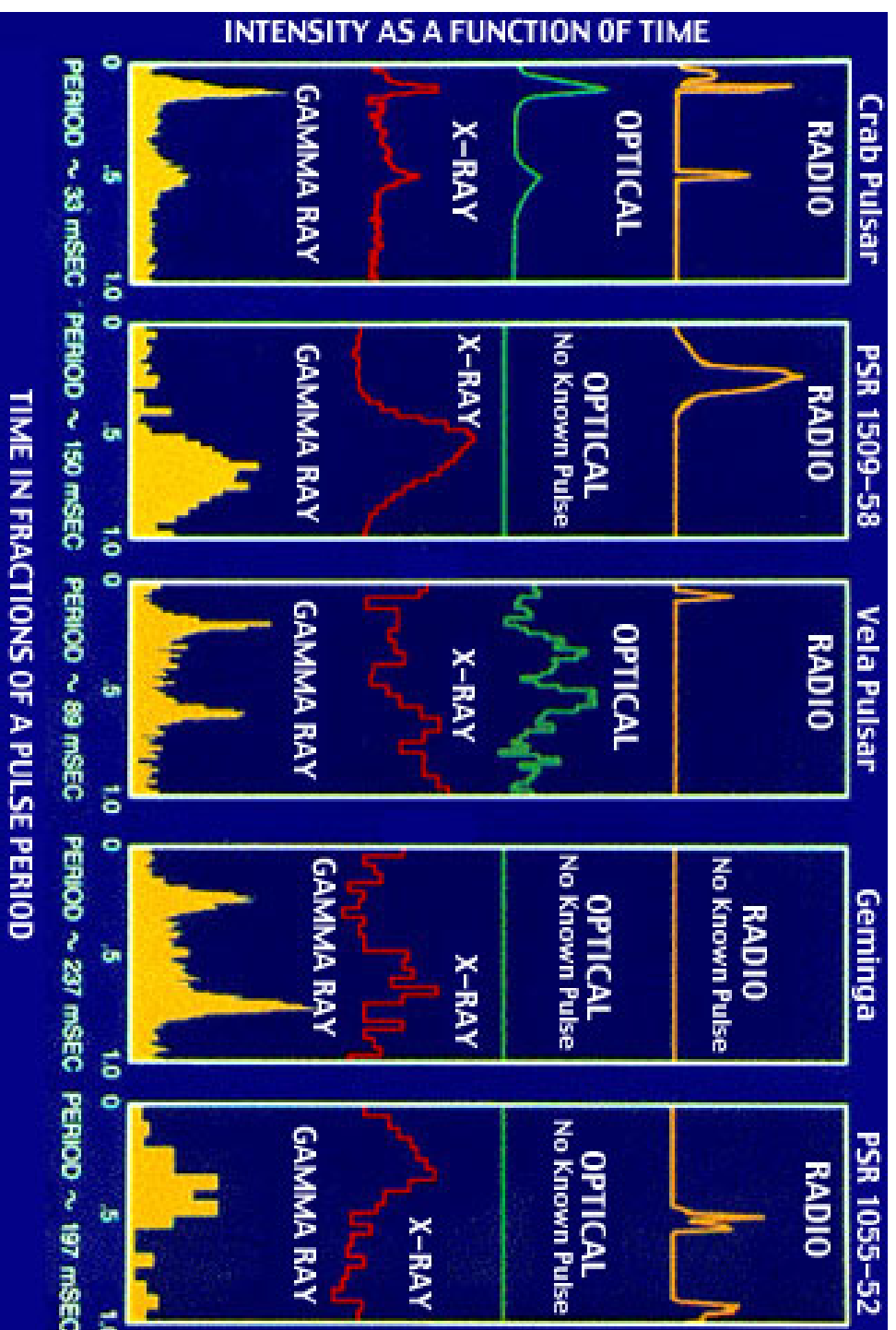
- values: $\dot{E} \sim (10^2 - 10^6) L_{\odot}$, brightest stars: $L \sim 10^6 L_{\odot}$
- total rotational energy of a young pulsar: $E_{rot} \sim 10^{51} \text{ erg}$

Pulsars are powerful galactic emission sources,

(young neutron stars cool by thermal neutrino and photon emission, E_{int} decreases)

b) Nature of pulsar emission

- detected in a broad range of electromagnetic radiation:
radio, optical, X-rays, γ -rays



- however total electromagnetic small: $L_{em} \ll \dot{E}$
e.g. Crab: $L_X + L_\gamma \sim 0.1 \dot{E}$, $L_{radio} \sim 10^{-8} \dot{E}$
("by-product" of pulsar's main activity)

2. Crab nebula

a) Main emission:

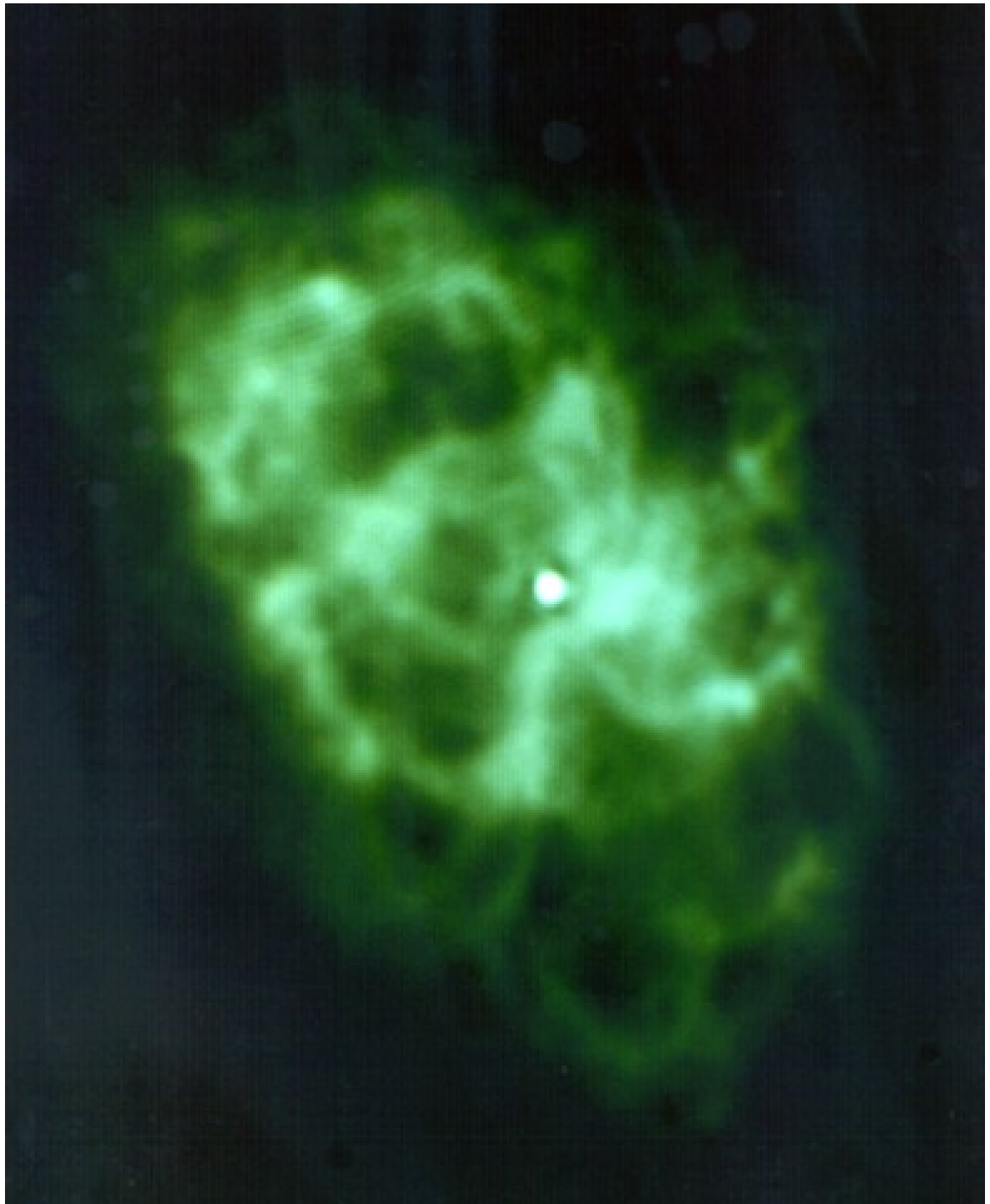
Bulk of \dot{E} emitted as a relativistic magnetized wind.

The pulsar wind visible when interacts with external matter.

b) Pulsar wind nebulae

- very young pulsars (~ 1000 years) still in SNR:
pulsar wind confined by external pressure

- Crab Nebula - a prototype (SN 1054):
expanding bubble, seen radio, optical and X-ray nebula

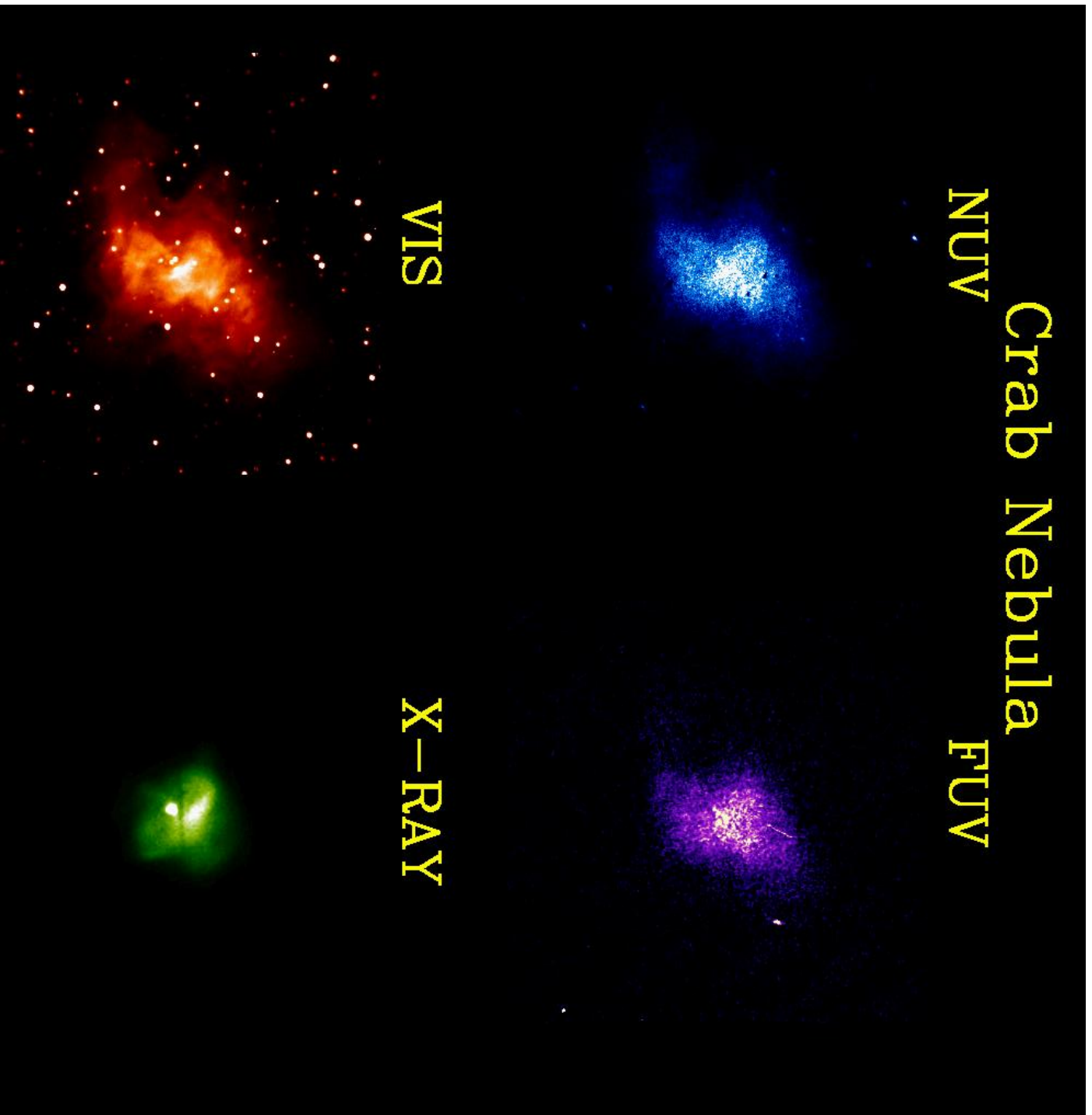


Crab Nebula

NUV FUV

VIS

X-RAY



- Synchrotron radiation by relativistic electrons (positrons) from the central pulsar

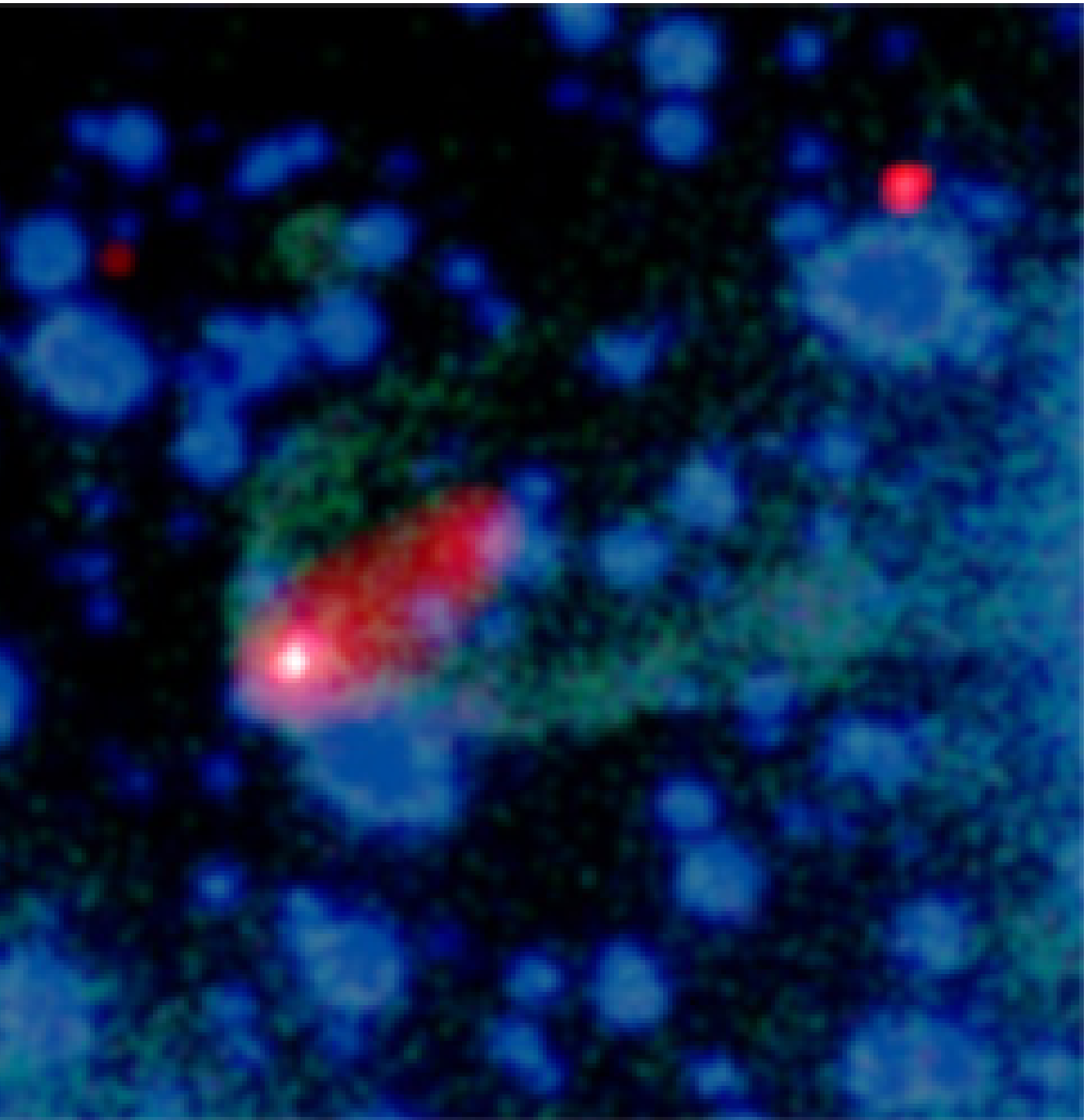
- radio emission - traces "integrated history" (long electron lifetime)

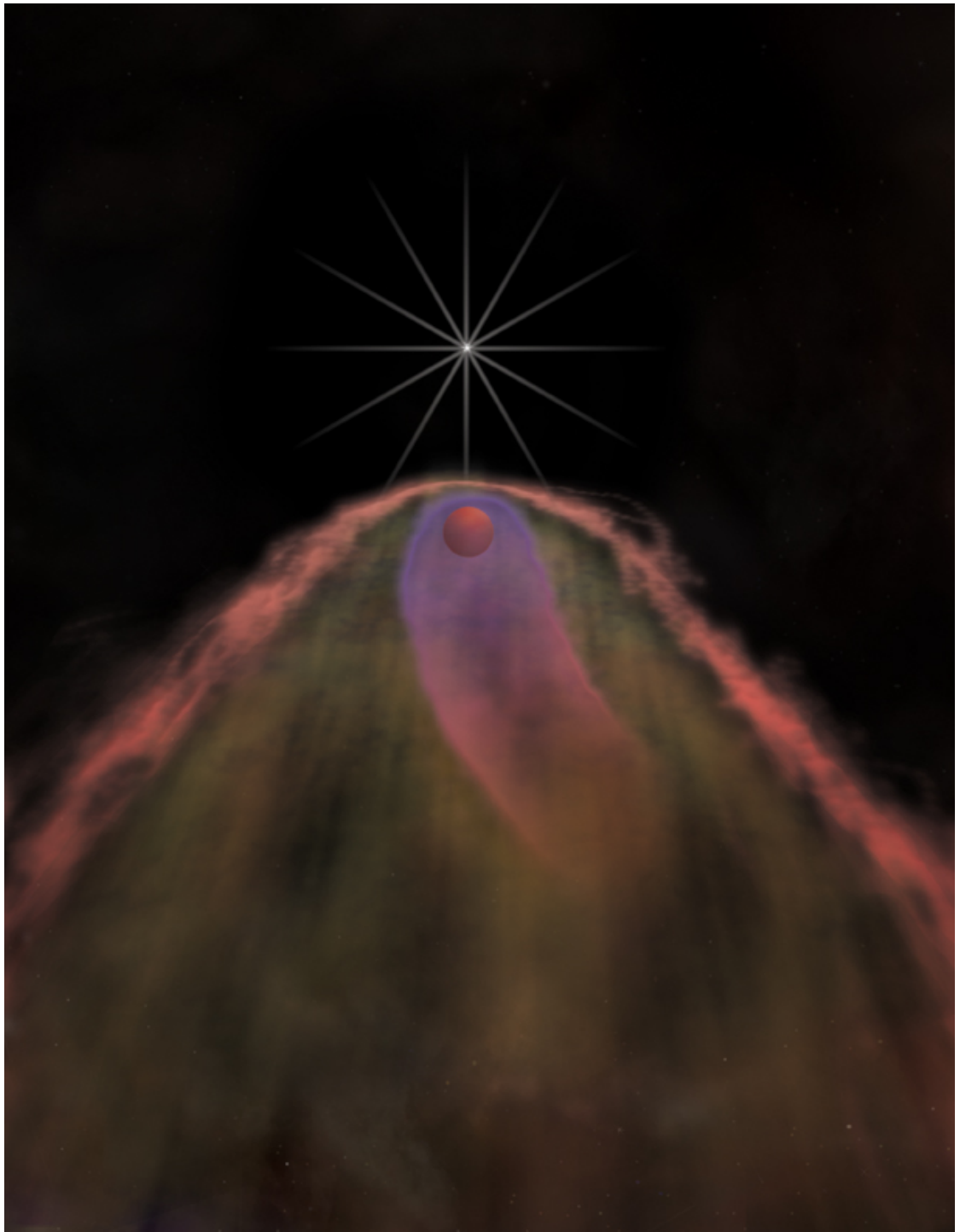
- X-ray emission - current pulsar state (short lifetime)

c) Bow-shock nebulae

Older pulsars: SNR dissipated, neutron stars move at $\sim 500 \text{ km s}^{-1}$.

Pulsar's wind generates a bow shock.





d) Properties of the Crab Nebula

- wind terminated by shock at $r_s = 0.1 pc$
- wind energy

$$\dot{E} = W_{Poynting} + W_{kinetic}$$

- model of Kennel and Coroniti (1984): very good fit of spectrum with $\gamma_{wind} = 3 \times 10^6$ and

$$\sigma = \frac{W_{Poynting}}{W_{kinetic}} = 0.003$$

before the shock.

- the wind at $r_s = 0.1 pc$ is dominated by kinetic energy of plasma with $\gamma \sim 10^6$

3. Pulsar wind models

a) Basic idea:

- rotating magnetized neutron star generates electromagnetic radiation (Pacini 1967, Gold 1968)
- strong electric field is induced, $\mathbf{E} = c^{-1} \boldsymbol{\Omega} \times \mathbf{r} \times \mathbf{B}$, which injects charged particles, e^+e^- pairs, into magnetosphere (Goldreich and Julian 1969; Strurrock 1970)
- electromagnetic field accelerates the plasma converting its energy into kinetic energy of the outflowing plasma (Poynting flux acceleration)

b) Historical model

Gunn and Ostriker (1969) (not realistic):

- magnetic dipole rotating with frequency $\Omega \rightarrow$
- electromagnetic monochromatic spherical waves of low frequency $\Omega \rightarrow$
- (test) particles tightly coupled to the wave by its strong magnetic field, coupling $\sim eB(mc\Omega)^{-1} > 10^8 \rightarrow$
- particles "ride the wave" at essentially constant phase;
- very effective acceleration up to the energy

$$E_c = mc^2 \left[\frac{3}{\sqrt{2}} \frac{eB_0}{mc\Omega} \ln \frac{r_c}{r_0} \right]^{\frac{2}{3}}$$

- Crab: $r_0 = c\Omega = 1576km$ (light cylinder),

$$r_c = r_s = 0.1pc$$

- protons: $E_c = 1.6 \times 10^{15} eV \approx 1.6 \times 10^6 m_p c^2$

unfortunately, pulsars do not radiate strong electromagnetic waves of rotation frequency Ω

c) Current (most promising) model

- observations hints: mechanism needed to effectively convert \dot{E}_{rot} into kinetic energy of the wind, employing magnetic field, with $\sigma \gg 1$ near the light cylinder and $\sigma \ll 1$ at $r \sim 0.1pc$.

-the problem of pulsar emission, especially converting

the Poynting flux into kinetic energy of the outflowing matter, turned out to be very complex and difficult (e.g. Michel 2001)

Recent results

- Contopoulos, Kazanas and Fendt (1999) first self-consistent solution of the axisymmetric magnetosphere of an aligned rotating magnetic dipole (MHD)
- split-monopole-like open field lines far from the light cylinder, $r \sin \theta \gg R_{lc}$
- field lines velocity

$$\mathbf{v}_E = c \frac{\mathbf{E} \times \mathbf{B}}{B^2} = c(\hat{r} \frac{x^2}{1+x^2} + \hat{\phi} \frac{x}{1+x^2})$$

$$x \equiv \frac{r \sin \theta}{R_{lc}}$$

$$\gamma_E = \left(1 - \left(\frac{v_E}{c}\right)^2\right)^{-\frac{1}{2}} \rightarrow \frac{r \sin \theta}{R_{lc}}$$

for $r \sin \theta \gg R_{lc}$:

- the wind Lorentz factor becomes $\gamma_{wind} \rightarrow \gamma_E$
- particles “surf-ride” on the electromagnetic field (B-parallel velocity negligible)

- the wind is a **linear** accelerator, $\mathbf{v}_E \sim \hat{\mathbf{r}}$; radiation losses **negligible**

4. Neutron stars with ultra-strong magnetic fields

- a) Maximum energy of particles (from possible potential difference)

$$E_{max} = 3 \times 10^{22} \frac{\mu}{10^{33} cgs} \left(\frac{\Omega}{10^4 s^{-1}} \right)^2 eV$$

- neutron star magnetic moment $\mu > 10^{32} cgs$ for

$$B > 10^{14} G:$$

$$B = 10^{12} G, P = 0.033s, E_{max} = 6 \times 10^{18} eV$$

$$\text{Maximum energy} \sim 10^{21} eV$$

only for ultra-strong magnetic fields, $B \sim 10^{14} G$

and fast rotation ($\Omega = 10^4 s^{-1}$, $P = 0.63ms$)

Most promising sources: fast rotating neutron stars with ultra-strong magnetic fields

b) High-field neutron stars

- Magnetars in Soft Gamma Repeters

(4 known): $B \sim 10^{14}G$,

periods now $\sim 5 - 10s$;

probably born with short periods $\sim 1ms$

- Anomalous X-ray Pulsars (AXP) similar B and P as magnetars (possibly the same class of neutron stars),
6 known (and candidates)

- Radio pulsars with magnetar field (recent discovery in Parkes Pulsar Survey) magnetic field $B = 10^{14}G$ and periods $P=3s, 6s$.